



Analytical Approximation for Cahn-Hilliard Phase-Field Model for Spinodal Decomposition of a Binary System

Ali Tozar¹ , Orkun Taşbozan² , Ali Kurt³ 

Article History

Received: 02 Oct 2020

Accepted: 23 Sep 2021

Published: 30 Sep 2021

10.53570/jnt.804302

Research Article

Abstract — Phase transformations which lead to dramatical property change are very important for engineering materials. Phase-field methods are one of the most successful and practical methods for modelling phase transformations in materials. The Cahn-Hilliard phase-field model is among the most promising phase-field models. The most successful aspect of the model is that it can predict spinodal decomposition (which is essential to determining the microstructure of an alloy) in a binary system. It is used in both materials science and many other fields, such as polymer science, astrophysics, and computer science. In this study, the Cahn-Hilliard phase-field model is evaluated by an analytical approach using the $(1/G')$ -expansion method. The solutions obtained are tested for certain thermodynamic conditions, and their accuracy of predicting the spinodal decomposition of a binary system is confirmed.

Keywords — Cahn-Hilliard phase field model, spinodal decomposition, $(1/G')$ -expansion method

Mathematics Subject Classification (2020) — 35R11, 34A08

1. Introduction

Foreseeing phase transformations is very important for tailoring the properties of engineering materials. For instance, occurrence of domains may determine magnetic properties of a metallic material [1]. On the other hand, mechanical properties of a material can be altered by a phase transformation [2]. Because of having a significant impact on the physical and chemical properties of the material, phase transformations, which are changes in the microstructure at equilibrium or non-equilibrium states, are the primary and most important concern of a materials engineering scientist. There are certain difficulties that this issue brings along with its importance. A method developed to model phase transformations to predict the properties of the material needs to take into account many thermal, electromagnetic, elastic and chemical effects. Hence, the issue is getting very difficult mathematically.

Phase-field methods are one of the most successful and practical methods for modeling phase transformations in materials. A phase field model can be described as spatiotemporal treating of the microstructure of a material as a scalar field (concentration, strain temperature, etc.). Phase field models have attracted much attention in recent years due to their versatility [3–5]. In these models, the free energy equation based on the thermodynamic state of the system is given in terms of phase fields. Nonlinear partial differential equations play an important role in the phase-field model due to dynamical governing of one or more order parameters as well as heat or mass transfer. These

¹tozarali@gmail.com; ²otasbozan@mku.edu.tr; ³pau.dr.alikurt@gmail.com (Corresponding Author)

¹Department of Physics, Hatay Mustafa Kemal University, Hatay, Turkey

²Department of Mathematics, Hatay Mustafa Kemal University, Hatay, Turkey

³Department of Mathematics, Pamukkale University, Denizli, Turkey

dynamical interactions of the parameters are based on the tendency of free energy to reach to its possible minimum value.

Cahn-Hilliard phase field model is the one of the most promising phase field models [6,7]. The most successful aspect of the model is that it can predict spinodal decomposition (which is very essential for determining the microstructure of an alloy) in a binary system [8]. Due to the success of the model in this regard, it is being used both in materials science [9,10] and in many fields such as polymer science [11], astrophysics [12], computer science [13] and astrophysics [12]. The simple structure of the model is mainly due to the fact that it is based on the diffusion laws in the Ginzburg-Landau theorem [14,15]. Total energy of a binary alloy is illustrated by the sum of two parts which are the bulk part depending on local energy difference and the interfacial part depending on concentration gradient. In summary, the total energy expression can be written as follows:

$$E = \int_{\Omega} f(\varphi) d\Omega + \int_{\Omega} \frac{\epsilon^2}{2} |\nabla \varphi|^2 d\Omega \quad (1)$$

where the first term is bulk and second term is interfacial energy ($\epsilon > 0$ represents the interfacial parameter). In order to evaluate the thermodynamic conditions, on the $f(\varphi)$ which is the free energy function is considered as follows according to Flory-Huggins model [16];

$$f(\varphi) = \frac{1}{2}((1 + \varphi) \ln(1 + \varphi) + (1 - \varphi) \ln(1 - \varphi) - \theta \varphi^2)$$

where $\theta = T/T_{critical}$ and φ is the phase field.

In the literature, different numerical techniques such as radial basis functions differential quadrature [17], finite element [18] and reduced differential transform [19] methods have been employed for the solution of Cahn-Hilliard equation. Unfortunately, numerical methods give approximate results. Analytical methods can give precise predictions to phase separation processes. Thus, Cahn-Hilliard phase field model is evaluated by an analytical approach using the $(1/G')$ -expansion method in this study. The obtained analytical solutions are firstly appeared in the literature.

2. Governing Equation

Authors considered the general form of the Cahn-Hilliard [20];

$$\varphi_t + \varphi_{xxxx} = (A(\varphi))_{xx} + l\varphi_x \cdot l > 0 \quad (2)$$

In this equation $A(\varphi(x, t))$ denotes an inherent chemical potential arising from driven forces (such as thermodynamic) that has the form of $A(\varphi(x, t)) = \varphi^3(x, t) - \varphi(x, t)$ and $\varphi(x, t)$ states phase field of a binary system that the components are having separated and $l\varphi_x$ indicates the diffusion induced kinetics of the mixture (such as a binary alloy).

3. $(1/G')$ -Expansion Method

Many authors used The $(1/G')$ -expansion method as a tool to get the exact solutions of various differential equations [21–23]. These studies represent that considered method is a powerful and effective while obtaining the analytical solutions of partial differential equations (PDEs). Suppose that nonlinear PDE for $\varphi(x, t)$ is given in the form

$$H \left(\varphi, \frac{\partial \varphi}{\partial t}, \frac{\partial \varphi}{\partial x}, \frac{\partial^2 \varphi}{\partial t^2}, \frac{\partial^2 \varphi}{\partial x^2}, \dots \right) = 0 \quad (3)$$

that the unknown function $\varphi(x, t)$ is the function of two independent the variables x, t and also H is the polynomial of the function $\varphi(x, t)$ and its higher order partial derivatives.

Regarding the wave variable as

$$\varphi(x, t) = \varphi(\xi), \xi = kx + ct \quad (4)$$

where k and c can be described as free constants. By using Equation (4) and Equation (3) changes into an nonlinear ordinary differential equation (ODE) with respect to $\varphi = \varphi(\xi)$

$$F(\varphi, \varphi', \varphi'', \varphi''', \dots) = 0 \tag{5}$$

where prime indicates Newtonian concept derivative due to ξ . According to considered method the traveling wave solutions of Equation (5) can be expressed as a polynomial of $(1/G')$ as follows

$$\varphi(\xi) = \sum_{i=0}^n a_i \left(\frac{1}{G'}\right)^i, \quad a_n \neq 0 \tag{6}$$

where $G = G(\xi)$ satisfies the following differential equation

$$G'' + \lambda G' + \mu = 0 \tag{7}$$

and $a_i (i = 0, \dots, n), \lambda, \mu$ are random constants to be examined later. To represent the solution of Equation (7) with $G = G(\xi)$, the Equation (6) will include the following equation

$$\frac{1}{G'(\xi)} = \frac{1}{-\frac{\mu}{\lambda} + A \tanh(\lambda\xi) - A \sinh(\lambda\xi)} \tag{8}$$

where A is integral constant.

Step 1. By using the homogeneous balance principle between the highest nonlinear terms and the highest order derivatives of $\varphi(\xi)$ in Equation (5), the positive integer n in Equation (6) can be determined.

Step 2. Subrogating (6) with Equation (7) into Equation (5) and collecting together all the same powered terms of $(1/G')$ together, the left hand side of Equation (5) is turns into a polynomial with respect to $(1/G')$. After equating each coefficient of this polynomial to zero, we handle an algebraic equation system with respect to $a_i (i = 0, \dots, n), \lambda, \mu, c, k$.

Step 3. In this step symbolic computer software is used to solve the algebraic equations system with respect to arbitrary constants $a_i (i = 0, \dots, n), \lambda, \mu, c, k$, then subrogating the results with the solutions of Equation (7) into Equation (6) led to traveling wave solutions of Equation (5).

4. The Analytical Solution of Cahn-Hilliard Equation

The transformation

$$\varphi(x, t) = \varphi(\xi), \xi = kx + ct \tag{9}$$

converts Equation (2) into nonlinear differential equation as

$$(c - l)\varphi + \varphi''' - (\varphi^3 - \varphi)' = 0 \tag{10}$$

Now, considering the homogeneous balance principle between φ''' and $\varphi^2\varphi'$ appearing in Equation (10), we get $n = 1$. Consequently, we can write the equation:

$$\varphi(\xi) = a_1 \left(\frac{1}{G'(\xi)}\right) + a_0, \quad a_1 \neq 0 \tag{11}$$

and therefore

$$\varphi'(\xi) = a_1\mu \left(\frac{1}{G'(\xi)}\right)^2 + a_1\lambda \left(\frac{1}{G'(\xi)}\right) \tag{12}$$

$$\varphi'''(\xi) = 6a_1\mu^3 \left(\frac{1}{G'(\xi)}\right)^4 + 12a_1\mu^2\lambda \left(\frac{1}{G'(\xi)}\right)^3 + 7a_1\lambda^2\mu \left(\frac{1}{G'(\xi)}\right)^2 + a_1\lambda^3 \left(\frac{1}{G'(\xi)}\right) \tag{13}$$

By substituting Equations (11)-(13) into Equation (10) and bringing all terms with the same power of $(1/G')$ together, the left-hand side of Equation (10) is changes into another polynomial in $(1/G')$.

Equating each coefficient of this polynomial to zero, yields a set of simultaneous algebraic equations for a_0, a_1, k, c, λ and μ follows:

$$\begin{aligned} \left(\frac{1}{G'(\xi)}\right)^0 &: ca_0 - lka_0 = 0 \\ \left(\frac{1}{G'(\xi)}\right)^1 &: ca_1 - lka_1 - 3k^2a_1a_0^2\lambda + k^2a_1\lambda + k^4a_1\lambda^3 = 0 \\ \left(\frac{1}{G'(\xi)}\right)^2 &: -3k^2a_1a_0^2\mu + 7k^4a_1\lambda^2\mu + k^2a_1\mu - 6k^2a_1^2a_0\lambda = 0 \\ \left(\frac{1}{G'(\xi)}\right)^3 &: -3k^2a_1^3\lambda + 12k^4a_1\mu^2\lambda - 6k^2a_1^2a_0\mu = 0 \\ \left(\frac{1}{G'(\xi)}\right)^4 &: 6k^4a_1\mu^3 - 3k^2a_1^3\mu = 0 \end{aligned}$$

Solving the algebraic equations above, yields

$$c = \pm \frac{l\sqrt{2}}{\lambda}, \quad k = \pm \frac{\sqrt{2}}{\lambda}, \quad a_0 = \pm 1, \quad a_1 = \pm \frac{2\mu}{\lambda} \tag{14}$$

and λ, μ are arbitrary constants. By using (14), expression (11) and (8), we have of travelling wave solutions of Equation (2) as follows:

$$\varphi_{1,2}(x, t) = \pm 1 \pm \frac{2\mu}{-\mu + A\lambda \cosh(\sqrt{2}lt + \sqrt{2}x) \mp A\lambda \sinh(\sqrt{2}lt + \sqrt{2}x)} \tag{15}$$

5. Graphical Representation of $\varphi_1(x, t)$ for Some Thermodynamic Conditions

Interfacial parameter ϵ thickening with the decrease in the l parameter can be seen in Figure 1(a). It can be said that in a binary system interfacial region where the metastable phase exist narrows with the increasing l values. As it is expected from the model that the double-well structure for the θ values bigger than one can be seen in the Figure 1(b). The peak between the troughs of the wells corresponds to the metastable phase in the interfacial region. This issue can be more clearly in Figure 2(a-c) which show the energy density distribution for the corresponding θ values of 1, 1.25, and 1.5.

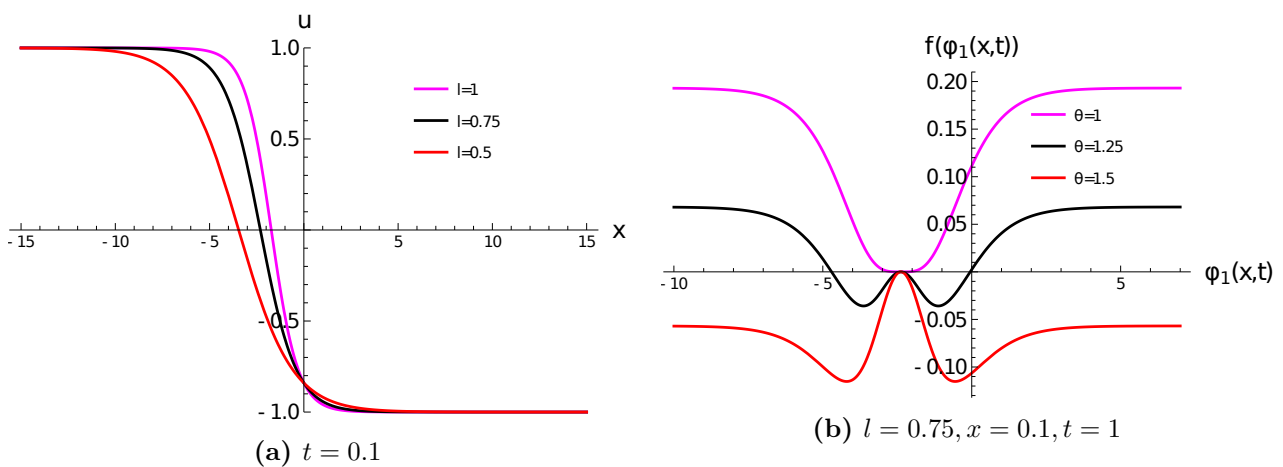


Fig. 1. Phase field (a), free energy (b) for the solution $\varphi_1(x, t)$ for $\mu = -0.1, \lambda = 0.1, A = 0.1$.

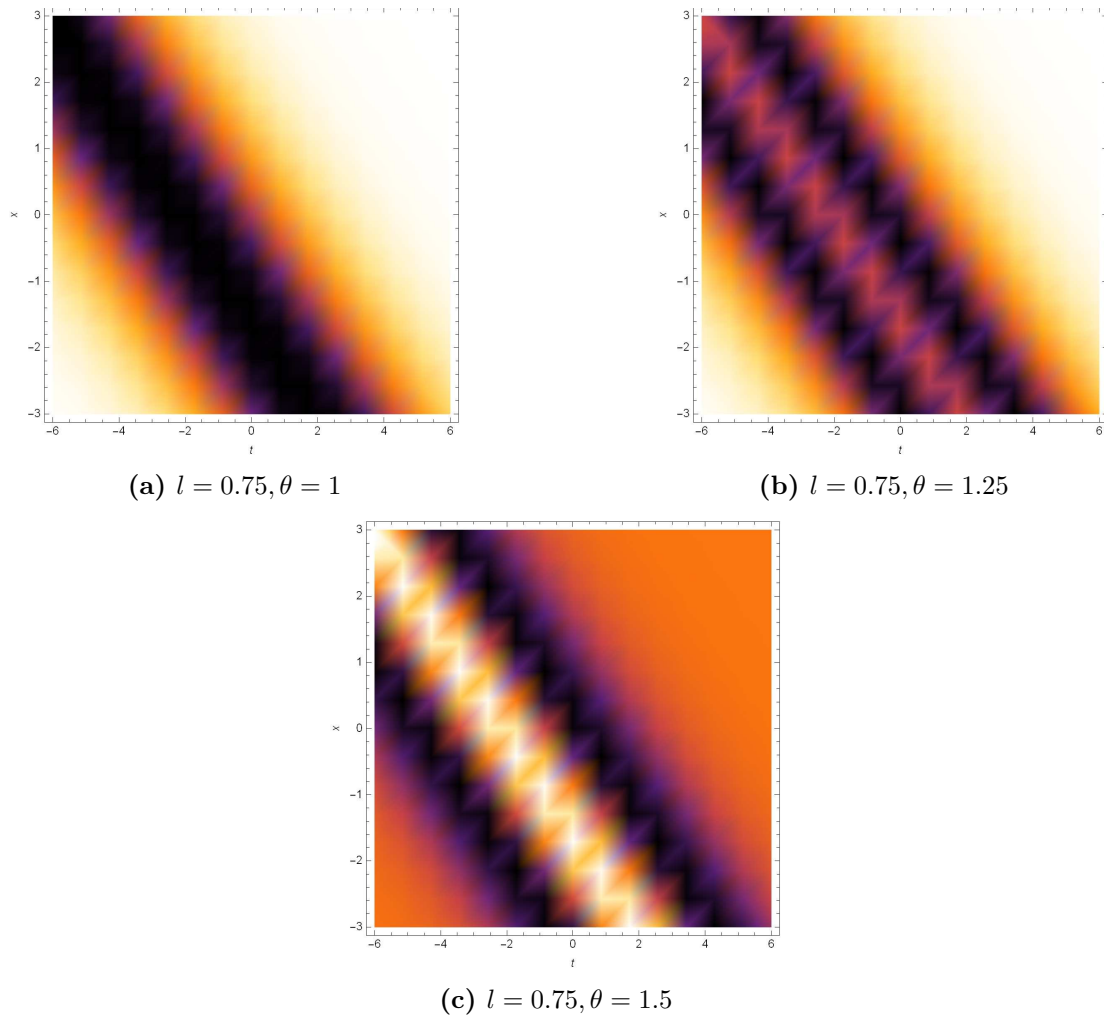


Fig. 2. Spatiotemporal energy distribution (a-c) for the solution $\varphi_1(x, t)$ for $\mu = -0.1, \lambda = 0.1, A = 0.1$.

6. Conclusion

The exact solutions of Cahn-Hillard model have been achieved by the $(1/G')$ -expansion method. An energy equation based on Flory-Huggins model was employed for evaluation of the solutions in certain thermodynamic conditions. As it is expected from the model double-well structure of the energy density distribution which corresponds to spinodal decomposition and occurrence of a metastable phase was seen. The results obtained in this study were shown to be accurate in predicting spinodal decomposition phenomenon. Also, some 2D graphical representations and contour plots of the obtained solutions are given for different values of l and θ . These results may be used by material engineers for tailoring the materials properties.

Author Contributions

All the authors contributed equally to this work. They all read and approved the last version of the manuscript.

Conflicts of Interest

The authors declare no conflict of interest.

References

- [1] V. P. Menushenkov, M. V. Gorshenkov, I. V. Shchetinin, A. G. Savchenko, E. S. Savchenko, D. G. Zhukov, *Evolution of the Microstructure and Magnetic Properties of As-Cast and Melt Spun Fe₂NiAl Alloy During Aging*, Journal of Magnetism and Magnetic Materials 390 (2015) 40–49.
- [2] H. R. Sistla, J. W. Newkirk, F. F. Liou, *Effect of Al/Ni ratio, Heat Treatment on Phase Transformations and Microstructure of Al_xFeCoCrNi_{2-x} (x=0.3, 1) High Entropy Alloys*, Materials Design 81 (2015) 113–121.
- [3] M. Javanbakht, V. I. Levitas, *Interaction between Phase Transformations and Dislocations at the Nanoscale. Part 2: Phase Field Simulation Examples*, Journal of the Mechanics and Physics of Solids 82 (2015) 164–185.
- [4] P. Steinmetz, Y. C. Yabansu, J. Hötzer, M. Jainta, B. Nestler, S. R. Kalidindi, *Analytics for Microstructure Datasets Produced by Phase-Field Simulations*, Acta Materialia 103 (2016) 192–203.
- [5] J. Kundin, L. Mushongera, H. Emmerich, *Phase-Field Modeling of Microstructure Formation During Rapid Solidification in Inconel 718 superalloy*, Acta Materialia 95 (2015) 343–356.
- [6] J. W. Cahn, J. E. Hilliard, *Free Energy of a Nonuniform System. I. Interfacial Free Energy*, The Journal of Chemical Physics 28 (1958) 258–267.
- [7] J. W. Cahn, *Free Energy of a Nonuniform System. II. Thermodynamic Basis*, The Journal of Chemical Physics 30 (1959) 1121–1124.
- [8] X. Zhang, G. Shen, C. W. Li, J. F. Gu, *Analysis of Interface Migration and Isothermal Martensite Formation for Quenching and Partitioning Process in a Low-Carbon Steel by Phase Field Modeling*, Modelling and Simulation in Materials Science and Engineering 27(7) (2019) 075011.
- [9] P. P. Moskvina, S. I. Skuratovskiy, O. P. Kravchenko, G. V. Skyba, H. V. Shapovalov, *Spinodal Decomposition and Composition Modulation Effect at the Low-Temperature Synthesis of Ax₃B_{1-x}C₅ Semiconductor Solid Solutions*, Journal of Crystal Growth, 510 (2019) 40–46.
- [10] N. Kuwahara, H. Sato, K. Kubota, *Kinetics of Spinodal Decomposition in a Polymer Mixture*, Physical Review E 47 (1993) 1132–1138.
- [11] D. Jeong, S. Lee, Y. Choi, J. Kim, *Energy-Minimizing Wavelengths of Equilibrium States for Diblock Copolymers in the Hex-Cylinder Phase*, Current Applied Physics 15 (2015) 799–804.
- [12] Y. F. Wang, Z. H. Xiao, S. Q. Shi, *Xe Gas Bubbles Evolution in UO₂ Fuels-A Phase Field Simulation*, Scientia Sinica: Physica, Mechanica et Astronomica 49 (2019) 11.
- [13] A. L. Bertozzi, S. Esedoglu, A. Gillette, *Inpainting of Binary Images Using the Cahn-Hilliard Equation*, IEEE Transactions on Image Processing 16 (2007) 285–291.
- [14] A. Vorobev, T. Lyubimova, *Vibrational Convection in a Heterogeneous Binary Mixture. Part 1. Time-Averaged Equations*, Journal of Fluid Mechanics 870 (2019) 543–562.
- [15] E. V. Radkevich, E. A. Lukashev, O. A. Vasil'eva, *Hydrodynamic Instabilities and Nonequilibrium Phase Transitions*, Doklady Mathematics 99 (2019) 308–312.
- [16] O. Wodo, B. Ganapathysubramanian, *Computationally Efficient Solution to the Cahn-Hilliard Equation: Adaptive Implicit Time Schemes, Mesh Sensitivity Analysis and the 3D Isoperimetric Problem*, Journal of Computational Physics 230 (2011) 6037–6060.

- [17] C. Liu, F. Frank, B. M. Rivière, *Numerical Error Analysis for Nonsymmetric Interior Penalty Discontinuous Galerkin Method of Cahn-Hilliard Equation*, Numerical Methods for Partial Differential Equations 35 (2019) 1509–1537.
- [18] M. Dehghan, V. Mohammadi, *The Numerical Solution of Cahn-Hilliard (CH) Equation in One, Two and Three-Dimensions via Globally Radial Basis Functions (GRBFs) and RBFs-Differential Quadrature (RBFs-DQ) Methods*, Engineering Analysis with Boundary Elements 51 (2015) 74–100.
- [19] A. M. S. Mahdy, N. A. H. Mukhtar, *Numerical Solution of Cahn-Hilliard Equation*, International Journal of Applied Engineering Research 13 (2018) 3150–3156.
- [20] D. Lu, M. S. Osman, M. M. A. Khater, R. A. M. Attia, D. Baleanu, *Analytical and Numerical Simulations for the Kinetics of Phase Separation in Iron (Fe-Cr-X (X=Mo, Cu)) based on Ternary Alloys*, Physica A: Statistical Mechanics and its Applications 537 (2020) 122634.
- [21] A. Yokuş, *An Expansion Method for Finding Traveling Wave Solutions to Nonlinear Pdes*, Istanbul Commerce University Journal of Science 14 (2015) 65–81.
- [22] H. Durur, A. Yokuş, *Hyperbolic Traveling Wave Solutions for Sawada Kotera Equation Using $(1/G')$ -Expansion Method*, Afyon Kocatepe University Journal of Sciences and Engineering 19 (2019) 615–619.
- [23] A. Yokuş, H. Durur, *Complex Hyperbolic Traveling Wave Solutions of Kuramoto-Sivashinsky Equation Using $(1/G')$ Expansion Method for Nonlinear Dynamic Theory*, Journal of Balikesir University Institute of Science and Technology 21 (2019) 590–599.