

Generation of Coherent XUV Radiation in Gas Mixture of N₂ and H₂ using High Power Laser System

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Geliş / Received: 23/12/2020, Kabul / Accepted: 05/05/2021

Abstract

High harmonic generation is studied in gas species, namely pure H₂, pure N₂, and their mixture in 1:1 ratio. Observation of harmonic orders from 17H to 35H orders is well resolved. The harmonic spectrum produced from pure H₂ is obtained, and the harmonic signal is weaker compared to that produced from pure N₂ and their mixture. On the other hand, harmonic yield in the gas mixture is enhanced compared to that from pure H₂. The maximum photon energy that harmonic orders can reach is calculated with the well-known cutoff formula. The experiment and the calculation results are in close agreement. In the mixture of N₂:H₂, the higher harmonics are enhanced compared to the plateau region harmonics. The enhancement of 35H orders in pure N₂ or H₂-N₂ mixture is observed compared to pure H₂.

Keywords: Atomic and molecular physics, gas mixture, femtosecond laser, high harmonic generation, extreme ultraviolet (XUV)

Yüksek Güçlü Lazer Sistemi Kullanılarak N₂ ve H₂ Gaz Karışımında Koherent XUV Radyasyonunun Üretilmesi

Öz

Saf H₂, saf N₂ ve bunların 1:1 oranında karışımları yüksek harmonik üretiminde çalışılmıştır. 17H'den 35H'ye kadar olan harmonik mertebeler gözlemlendi. Saf H₂'den üretilen harmonik spektrum elde edildi, ancak harmonik sinyal, saf N₂ ve bunların karışımından üretilen harmonik spektruma kıyasla daha zayıftır. Gaz karışımındaki harmonik verim, saf H₂'den gelen harmoniklere kıyasla artar. Harmonik mertebelerin ulaşabileceği maksimum foton enerjisi, iyi bilinen cutoff formülü ile hesaplandı. Deney ve hesaplama sonuçlarıyla elde edilen harmonik mertebeler bir uyum içindedir. N₂:H₂ karışımında, plato bölgesi harmoniklerine kıyasla daha yüksek harmonikler geliştirilmiştir. Saf N₂ veya H₂-N₂ karışımında 35H mertebelerinde saf H₂'ye göre artış gözlemlendi.

Anahtar Kelimeler: Atomik ve moleküler fiziği, gaz karışımı, femtosaniye lazer, yüksek harmonik üretimi, aşırı ultraviyole (XUV).

1. Introduction

High-order harmonics generation (HHG) is a way to generate a coherent powerful extreme ultraviolet (XUV) radiation source (Chang and Corkum, 2010; Kapteyn et al., 2005). This coherent XUV source has potential for many practical applications, such as the study of the atomic and molecular process (Remetter et al., 2006), quantum dynamics of the electronic wave packet (Paul et al., 2001), biological imaging (Eisebitt et al., 2004; Itatani et al., 2004; Ravasio

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et al., 2009; Seaberg et al., 2011), production of isolated attosecond laser pulses (Hentschel et al., 2001; Lewenstein, 2002), and so on. The wide range of applications motivates the development of high harmonic generation. The improvement of HHG has opened the realization of extreme nonlinear optics (Popmintchev et al., 2010), quantum interference effect (Vozzi et al., 2006), and molecular self-probing (Haessler et al., 2011). However, the bottleneck of this powerful light source is its low conversion energy, i.e. conversion from fundamental field to its higher frequencies. There are many experimental and theoretical studies to increase the efficiency of HHG in pure atomic and molecular gases (Brizuela et al., 2013; Ishikawa, 2003; Kanai et al., 2007; Takahashi et al., 2007). However, there are still many open questions to optimize the efficiency of the high harmonics because HHG strongly depends on macroscopic conditions such as phase matching, which can vary for different experimental approaches and designs.

Microscopic physics of harmonic generation can be well explained with a three-step model (Corkum, 1993; Lewenstein et al., 1994). In this model, an electron exposed to an intense laser field tunnels through the atomic barrier, and then the electron is driven back and forth with the oscillating laser field. In this process, the electron gains kinetic energy. Finally, the electron could recombine with its core and releases its kinetic energy in the form of high-energy photons. The highest harmonic energy (cutoff) is given by $E_{\max} = I_p + 3.17U_p$ (Lewenstein et al. 1994), where I_p is the ionization potential of the used gas, $U_p \sim I\lambda^2$ is the ponderomotive energy gained by the free electron due to the laser oscillation, and it is proportional to the laser intensity (I) and wavelength (λ) (Corkum, 1993; Lewenstein et al. 1994). There are several parameters to control the efficiency of HHs and to extend the cutoff harmonics. For example, the highest harmonic energy is proportional to the laser intensity (I) and the wavelength (λ). By controlling these two variables, one can optimize the harmonic signal and can extend the cutoff harmonic. On the other hand, there is a limitation to increase laser intensity ($\sim 10^{17}$ W/cm²) because the magnetic component of the laser field is comparable to the electric field of the pulse.

In this paper, we report harmonic generation from pure N₂, pure H₂, and their mixture using the high power laser system producing optical pulses having energy per pulse 6mJ with 50fs pulse duration at a 10 Hz repetition rate. The highest harmonics from 31H to 35H in the mixed gas medium are enhanced as a factor of ~ 3 and ~ 4 according to H₂, (high conversion efficiency). The results could be explained as that the harmonics produced in N₂ medium boost harmonics produced in pure H₂. The effect of the gas mixture has the potential to increase or decrease the efficiency of the HHG. Ionization potential (I_p) of both H₂ and N₂ are close to each other (I_p of N₂ molecules is 15.7eV (Tate and Smith, 1932)) and I_p of H₂ molecule is 15.4eV (Bleakney, 1932)), and they present well-resolved harmonic peaks. Therefore, the H₂-N₂ gas mixture is chosen for investigating the characteristic of harmonic generation.

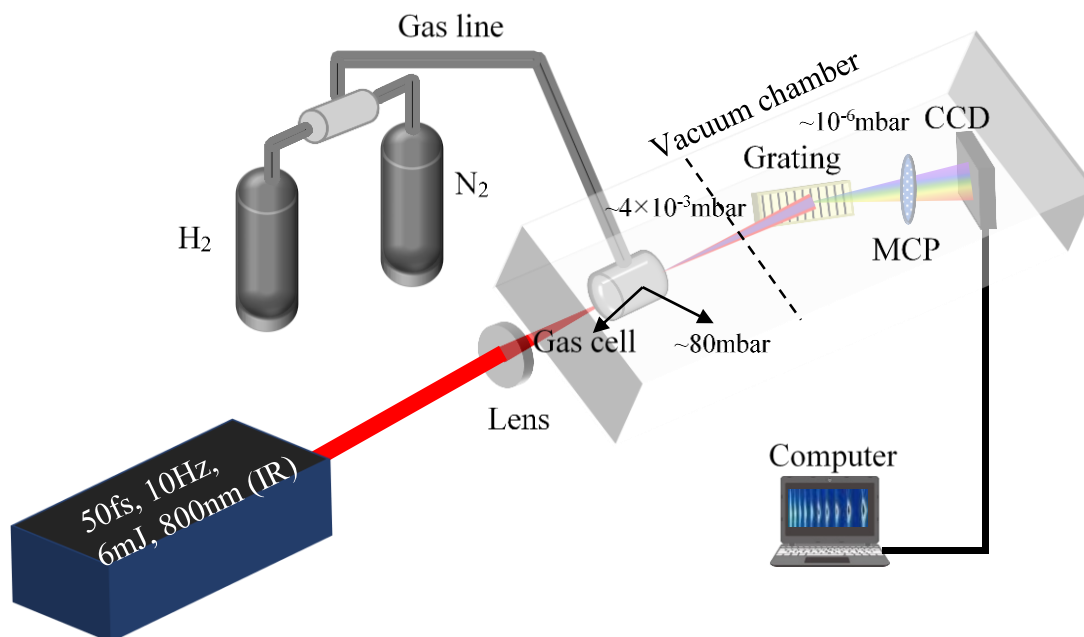


Figure 1. Schematic of the harmonic generation experimental setup. The setup includes a gas cell (harmonic source), a grating (for dispersion of harmonic spectrum), MCP (to amplify the harmonic yield), CCD (to capture the harmonics signal), IR (infrared radiation)

2. Experimental Setup

The laser system has been used that the driving laser field consists of high power laser system producing optical pulses about a 10Hz repetition rate, 50fs pulse duration (~19cycle) centered at 800nm, and energy of 6mJ per pulse. To generate high-order harmonics, the radiation is focused on a gas cell, which is machined in a machine shop to generate high harmonics. The gas cell has a length of 7mm and a diameter of 5mm. The cell has input and output holes, which are taped by aluminum foil tape to localize gas density inside the cell. The strong laser field is focused on the gas cell by using a 40cm focal length lens, and high-order harmonics of the fundamental field are generated. The intensity at the interaction region is estimated at $1 \times 10^{15} \text{Wcm}^{-2}$. After the gas cell, the generated harmonics propagate through to the spectrometer, and they are recorded. The schematic of the experimental design is presented in Fig. 1. The recorded data are analyzed by using the MATLAB program.

The powerful laser system allows one to control harmonic orders, and it can reach a high cutoff harmonics due to its high power. The generated harmonics propagate through the spectrometer. First, each harmonic order is reflected from the grating at a different angle. Second, the harmonics are amplified by a microchannel plate (MCP), and they are finally captured by the charge-coupled device (CCD) camera. The state-of-the-art laser system allows reaching higher-order harmonics by enhancing the cutoff harmonic. The cutoff harmonic in the mixed gases is decreased compared to that in pure N₂. The results can be explained as the harmonics coming from N₂ are absorbed by the H₂. Moreover, destructive interference of both species leads to this decrement. In addition, the ionization potential of both gases is close to each other, and so the phase mismatch induced by the free electron could cause this decrease in the mixture. However, the harmonic yield in the gas mixture is boosted compared to that produced in pure H₂. This mechanism could be constructive interference between harmonics produced in N₂ and H₂ gas species, respectively (Kanai et al., 2007).

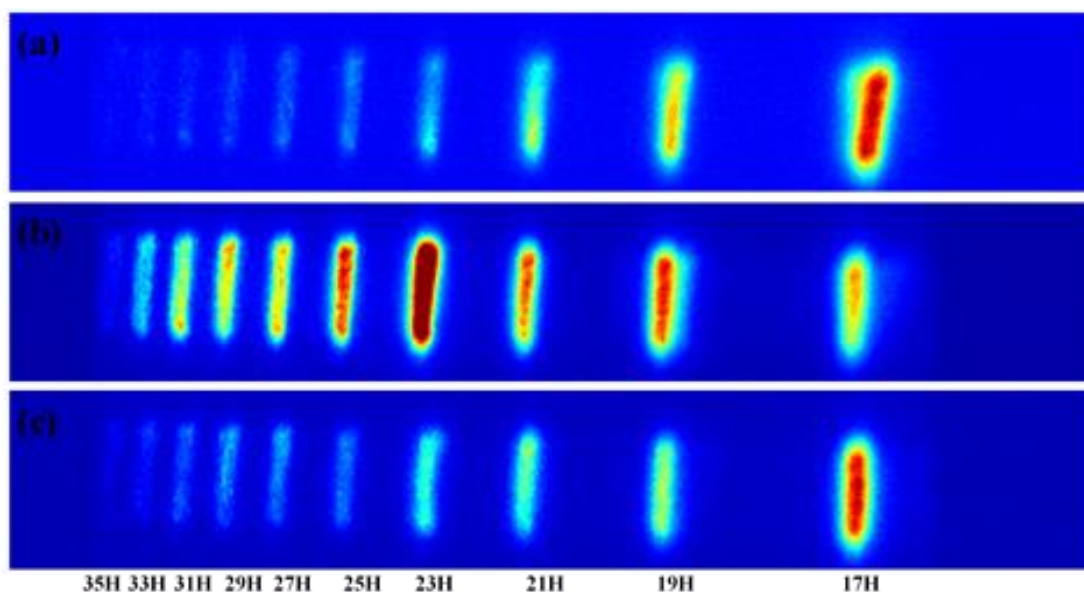


Figure 2. Captured Harmonic images by a CCD camera at the cell pressure of 80mbar. (a) pure H₂, (b) pure N₂ (c) N₂-H₂ mixture in 1:1 ratio

3. Experimental Results

Harmonic spectra are produced in three different gas mediums, N₂, H₂, and their mixture N₂-H₂, respectively. Figure 2 presents the typical CCD image of harmonic spectrum in (a) pure H₂, (b) pure N₂, and (c) N₂-H₂ mixture (50%-50% ratio) at the pressure of 80mbar. Harmonic orders are produced up to 35th harmonic orders, Fig. 2. The interaction chamber pressure is 4×10^{-3} mbar, which corresponds to interaction region pressure, ~ 80 mbar. The integration time of the CCD camera is set to 10s. The fundamental laser field power is 60mW corresponding $\sim 1 \times 10^{15}$ Wcm⁻² focused intensity.

The above-mentioned cutoff formula calculates the maximum photon energy of generated harmonics. The cutoff harmonic order is calculated as the 33rd harmonic order for N₂ and H₂ gas species at $\sim 1 \times 10^{15}$ Wcm⁻² IR laser intensity because the ionization potentials of these two species are close to each other. However, the harmonic signal in pure N₂ is stronger than that in pure H₂. This is attributed to electron dispersion and laser defocusing due to free electron distribution in the medium (Kanai et al., 2007). In addition, the generation of high harmonics in N₂-H₂ gas mixture produces a weak harmonic signal compared to the harmonic signal in pure N₂. This could be due to the absorption of harmonics (produced in N₂) by the H₂ molecules. There is another possibility that there is bright harmonics radiation since the ionization potentials of both species are close and lower. However, the low ionization potential of the medium results in serious electron dispersion and ionization-induced laser defocusing. These drawbacks result in a decrease of the harmonic signal in the gas mixture (Kanai et al., 2007).

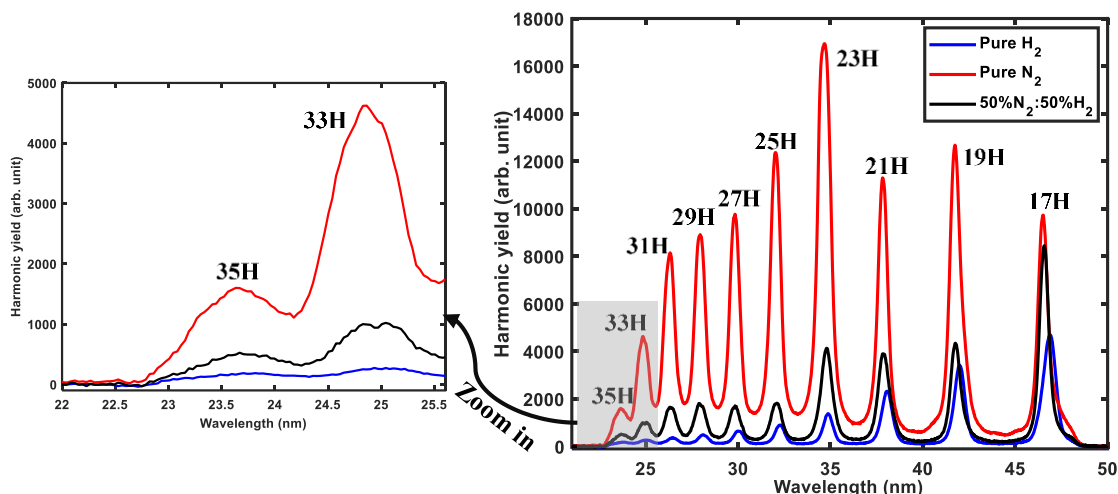


Figure 3. Harmonic spectrum in H₂, N₂, and their mixture with 50%:50% ratio. The solid blue line is for H₂. The solid red line is for N₂. The solid black line is for the mixture. Zoomed in part is for the harmonic spectrum for 33H and 35H orders

Figure 3 is for harmonic yield as a function of wavelength (nm) and photon energy (eV). The solid blue line is for H₂, the solid red line is for N₂, and the solid black line is for the gas mixture. The harmonics are well resolved up to 35 order in all gas species. The gas mixture produces a strong harmonic signal compared to the harmonic signal in pure H₂, Fig. 3. The generated strong harmonics in N₂ gas and the fundamental field (IR) lead to ionization of the H₂ gas in N₂-H₂ gas mixture. The mechanism can lead to an increase of the harmonic signal in N₂-H₂ gas mixture compared to the harmonic signal in pure H₂. Zoomed part of Figure 3 shows the 35th harmonic signal in pure N₂ and the mixture of N₂-H₂.

The decrease of the harmonic signal in the gas mixture of N₂-H₂ compared to the harmonic spectrum in pure N₂ can be explained because of the absorption of the harmonics in the gas mixtures. The absorption in the harmonic spectrum leads to a decrease of the harmonic signal. This paper is an experimental study to increase or decrease the high harmonic signal in the gas mixture. Figure 4 shows harmonic yield as a function of harmonic orders. Figures 4 (a) and (b) present a well harmonic distribution, plateau (19H-31H), and cutoff region (35H). Figure 4 (c) is for the ratio of harmonic yield in the mixture to harmonic yield in pure H₂. The harmonic yield in N₂-H₂ gas mixture is increased (about a factor of ~2 to 4) compared to the harmonic signal in pure H₂, Fig. 4(c). The 35H order produced in pure H₂ is relatively weak compared to other cases. In addition, harmonics near the cutoff region (31H to 35H) are more enhanced compared to those in the plateau region (19H-29H) for the H₂:N₂ gas mixture, Fig. 4(c)).

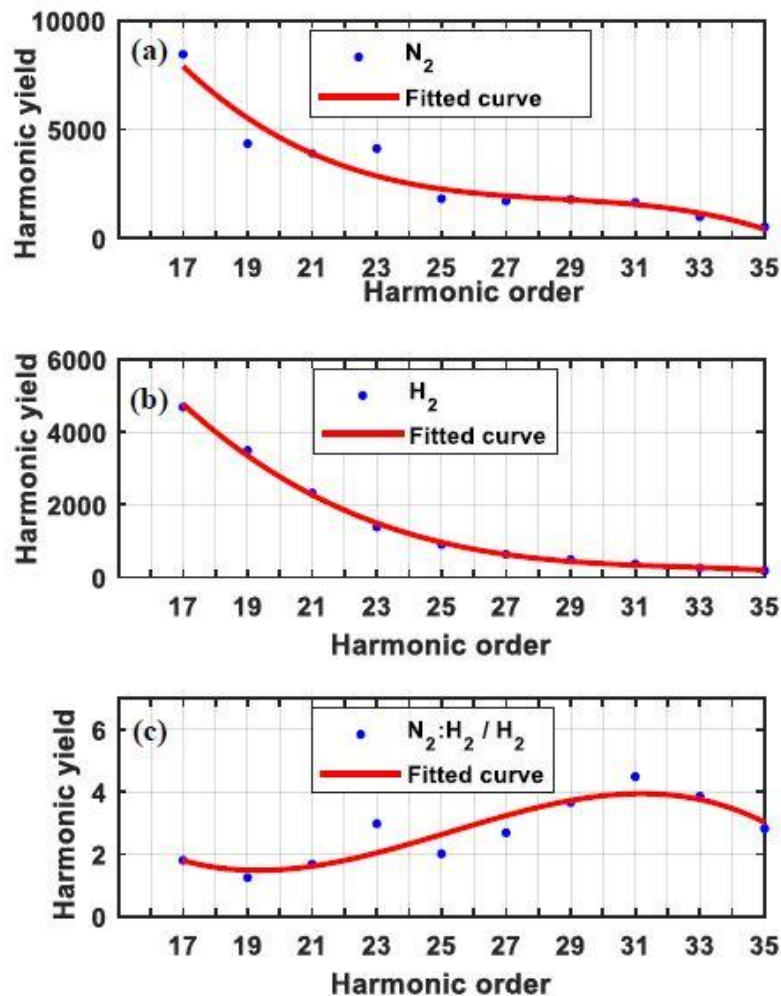


Figure 4. Harmonic yield as a function of harmonic orders. (a) N₂ harmonic yield (b) H₂ harmonic yield (c) N₂:H₂/H₂ harmonic yield. Blue dots are for experimental measurement, and the solid red line is for fitting

The curve fitting value for Fig. 4 is best performed by using a third-degree polynomial. The function of the fitting is $f(x) = P_1x^3 + P_2x^2 + P_3x + P_4$, where coefficients (P_1 , P_2 , P_3 , and P_4) are given in Table 1. Table 1 presents the harmonic yield polynomial fit for three different gas species. R-square takes values between 0 and 1, and it measures how successful the fit is performed. The value closer to 1 indicates the goodness of the fit. Polynomial coefficients and R square values can be seen in Table 1

Table 1. Fitting function coefficient for Figure 4.

Fitting Function: $f(x) = P_1x^3 + P_2x^2 + P_3x + P_4$					
Gas Species	P ₁	P ₂	P ₃	P ₄	R-square
N ₂	-719.2	940.1	-926.5	2084	0.9283
H ₂	-1.079	105.3	-3451	3.833×10^4	0.9977
(N ₂ : H ₂) / H ₂	-0.002944	0.2231	-5.323	42.28	0.802

4. Conclusion

Generation of high harmonics extending to ~54eV photon energy is experimentally produced in pure N₂, pure H₂, and their mixture in 1:1 ratio. The focused laser intensity is $1 \times 10^{15} \text{Wcm}^{-2}$ on the gas cell. A high harmonic spectrum of up to 35H orders is generated. The generation of the 35H is increased in the gas mixture compared to the harmonic signal produced in pure H₂ gas. Harmonic spectrums in N₂, H₂, and N₂-H₂ mixture are well resolved from 17H to 35H. The harmonic signal is increased in the gas mixture compared to harmonic yield in pure gas species of H₂ (Fig. 4). The mechanism of harmonic spectrum in the gas mixture can be explained that the harmonics produced in one gas species (N₂) with the fundamental field (IR) result in to increase in the harmonic signal from another gas species (H₂). The harmonic yield as a function of harmonic orders presents harmonic distribution, Fig. 4. Especially, higher orders are enhanced as a factor of 3 to 4. The fitting of harmonic yield is performed with R-squared value close to 1. The optimized harmonic source can be useful for various application areas such as attosecond light source, nonlinear optic in XUV region, and imaging of small-scale object by using harmonic source.

5. Acknowledgments

I am thankful to A. Kolomenski and H. Schuessler for providing the opportunity to perform the experiments described in this paper at Texas A&M University.

References

- Bleakney, W. (1932). "The ionization potential of molecular hydrogen", *Physical Review Letters*, 40, 496-501.
- Brizuela, F., Heyl, C. M., Rudawski, P., Kroon, D., Rading, L., Dahlström, J. M., Mauritsson, J., Johnsson, P., Arnold, C. L., and L'Huillier, A. (2013). "Efficient high-order harmonic generation boosted by below-threshold harmonics", *Scientific Reports*, 3, 1410.
- Chang, Z., and Corkum, P. (2010). "Attosecond photon sources: the first decade and beyond", *Journal of the Optical Society of America B*, 27, B9-B17.
- Corkum, P. B. (1993). "Plasma perspective on strong field multiphoton ionization", *Physical Review Letters*, 71, 1994-1997.
- Eisebitt, S., Luning, J., Schlotter, W. F., Lorgen, M., Hellwig, O., Eberhardt, W., and Stohr, J. (2004). "Lensless imaging of magnetic nanostructures by X-ray spectro-holography", *Nature*, 432, 885-888.
- Haessler, S., Caillat, J., and Salières, P. (2011). "Self-probing of molecules with high harmonic generation", *The Journal of Physics B: Atomic, Molecular and Optical Physics*, 44, 203001.
- Hentschel, M., Kienberger, R., Spielmann, C., Reider, G. A., Milosevic, N., Brabec, T., Corkum, P., Heinzmann, U., Drescher, M., and Krausz, F. (2001). "Attosecond metrology", *Nature*, 414, 509-513.
- Ishikawa, K. (2003). "Photoemission and ionization of He⁺ under simultaneous irradiation of fundamental laser and high-order harmonic pulses", *Physical Review Letters*, 91, 043002.

Itatani, J., Levesque, J., Zeidler, D., Niihara, H., Pepin, H., Kieffer, J. C., Corkum, P. B., and Villeneuve, D. M. (2004). "Tomographic imaging of molecular orbitals", *Nature*, 432, 867-871.

Kanai, T., Takahashi, E. J., Nabekawa, Y., and Midorikawa, K. (2007). "Destructive interference during high harmonic generation in mixed gases", *Physical Review Letters*, 98, 153904.

Kapteyn, H. C., Murnane, M. M., and Christov, I. P. (2005). "Extreme nonlinear optics: Coherent X rays from lasers", *Physics Today*, 58, 39-44.

Lewenstein, M. (2002). "Resolving physical processes on the attosecond time scale", *Science*, 297, 1131-1132.

Lewenstein, M., Balcou, P., Ivanov, M. Y., L'Huillier, A., and Corkum, P. B. (1994). "Theory of high-harmonic generation by low-frequency laser fields", *Physical Review A*, 49, 2117-2132.

Paul, P. M., Toma, E. S., Breger, P., Mullot, G., Augé, F., Balcou, P., Muller, H. G., and Agostini, P., (2001). "Observation of a train of attosecond pulses from high harmonic generation", *Science*, 292, 1689-1692.

Popmintchev, T., Chen, M. C., Arpin, P., Murnane, M. M., and Kapteyn, H. C. (2010). "The attosecond nonlinear optics of bright coherent X-ray generation", *Nature Photonics*, 4, 822-832.

Ravasio, A., Gauthier, D., Maia, F. R. N. C., Billon, M., Caumes, J. P., Garzella, D., Géléoc, M., Gobert, O., Hergott, J. F., Pena, A. M., Perez, H., Carré, B., Bourhis, E., Gierak, J., Madouri, A., Maily, D., Schiedt, B., Fajardo, M., Gautier, J., Zeitoun, P., Bucksbaum, P. H., Hajdu, J., and Merdji, H. (2009). "Single-Shot Diffractive Imaging with a Table-Top Femtosecond Soft X-Ray Laser-Harmonics Source", *Physical Review Letters*, 103, 028104.

Remetter, T., Johnsson, P., Mauritsson, J., Varju, K., Ni, Y., Lepine, F., Gustafsson, E., Kling, M., Khan, J., Lopez-Martens, R., Schafer, K. J., Vrakking, M. J. J., and L'Huillier, A. (2006). "Attosecond electron wave packet interferometry", *Nature Physics*, 2, 323-326.

Seaberg, M. D., Adams, D. E., Townsend, E. L., Raymondson, D. A., Schlotter, W. F., Liu, Y., Menoni, C. S., Rong, L., Chen, C.-C., Miao, J., Kapteyn, H. C., and Murnane, M. M. (2011). "Ultrahigh 22 nm resolution coherent diffractive imaging using a desktop 13 nm high harmonic source", *Optics Express*, 19, 22470-22479.

Takahashi, E. J., Kanai, T., Ishikawa, K. L., Nabekawa, Y., and Midorikawa, K. (2007). "Dramatic enhancement of high-order harmonic generation", *Physical Review Letters*, 99, 053904.

Tate, J. T. and Smith, P. T. (1932). "The Efficiencies of Ionization and Ionization Potentials of Various Gases Under Electron Impact", *Physical Review*, 39, 270-277.

Vozzi, C., Calegari, F., Benedetti, E., Berlasso, R., Sansone, G., Stagira, S., Nisoli, M., Altucci, C., Velotta, R., Torres, R., Heesel, E., Kajumba, N., and Marangos, J. P. (2006). "Probing two-centre interference in molecular high harmonic generation", *Journal of Physics B: Atomic, Molecular and Optical Physics*, 39, 457-466.