

THEORETICAL INVESTIGATION FOR YTTERBIUM EFFECT ON RADIATION SHIELDING CHARACTERISTICS OF 50Bi₂O₃-15Li₂O-15PbO-(20-X)B₂O₃-Yb₂O₃ BORATE GLASSES

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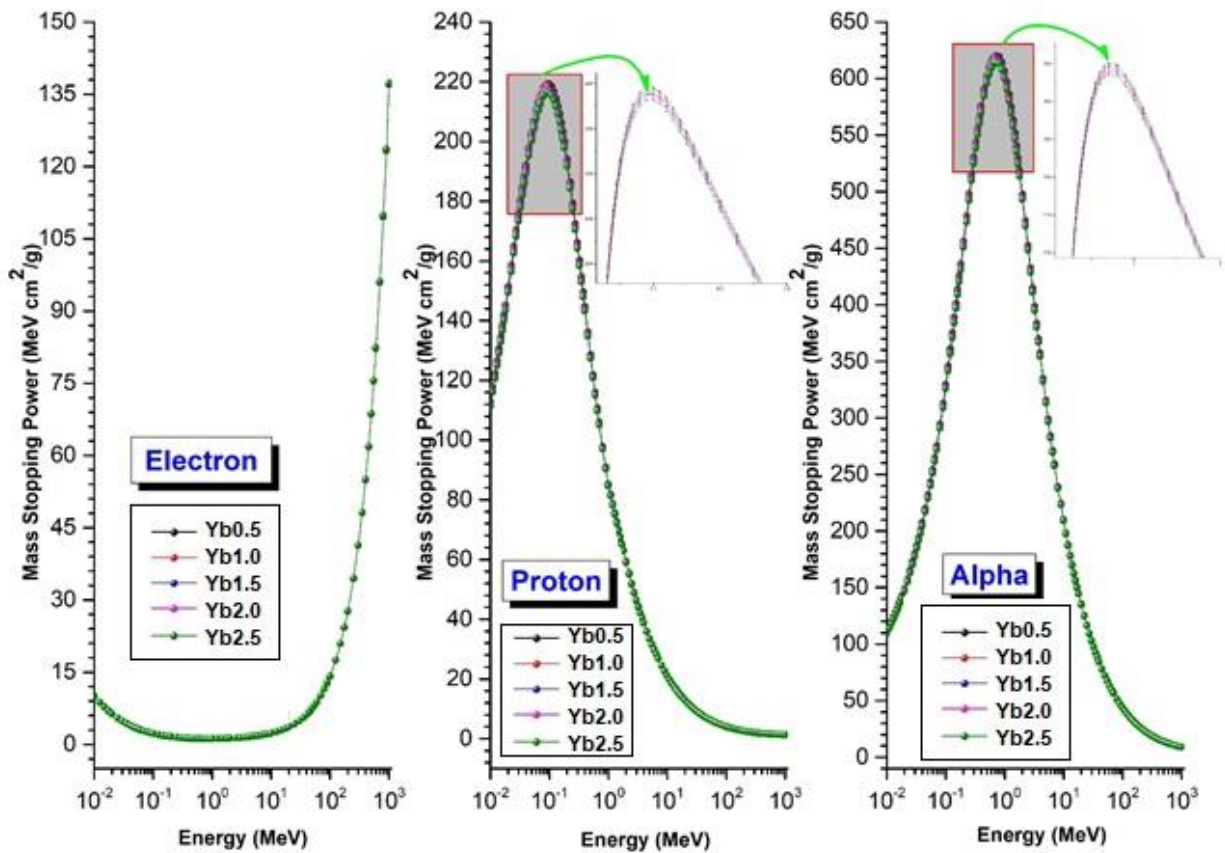
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Highlights

- Radiation shielding characteristics of ytterbium doped glasses have been studied.
- Phy-X/PSD, PAGEX, ESTAR, and SRIM codes have been used.
- Increasing amount of ytterbium exhibited higher shielding performances.

Graphical Abstract



Dependence of MSP values of alpha, proton and electron particles versus the energy.



THEORETICAL INVESTIGATION FOR YTTERBIUM EFFECT ON RADIATION SHIELDING CHARACTERISTICS OF $50\text{Bi}_2\text{O}_3\text{-}15\text{Li}_2\text{O-}15\text{PbO-(}20\text{-X)B}_2\text{O}_3\text{-Yb}_2\text{O}_3$ BORATE GLASSES

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ABSTRACT: As the nuclear industry has developed and radiation technologies have become more widely used, the dose of radiation from both synthetically produced sources and radioisotopes has increased, as has the number of people being irradiated. It is of the utmost importance to utilise appropriate shielding materials in order to reduce the negative effects of radiation sources. Doped glasses are among the most significant candidate materials in the field of radiation shielding. To this end, it was aimed to investigate the charged particle, gamma-ray and neutron shielding characteristics of ytterbium doped glasses with composition of $50\text{Bi}_2\text{O}_3\text{-}15\text{Li}_2\text{O-}15\text{PbO-(}20\text{-x)B}_2\text{O}_3\text{-Yb}_2\text{O}_3$ (where $x = 0.5, 1.0, 1.5, 2.0$ and 2.5). In this regard, the radiation shielding parameters were estimated by using Phy-X/PSD, PAGEX, ESTAR, and SRIM codes. Furthermore, the results were compared comprehensively and comparatively. Charged particle (alpha, proton and electron), gamma ray and neutron shielding efficiencies were found to be proportional to ytterbium content. The glasses with the increasing amount of ytterbium exhibited higher shielding performances. It can be said that all the studied glasses and particularly Yb2.5 sample could be used as shielding materials in many radiation related applications.

Keywords: Borate Glass, Charged Particle Shielding, Gamma Shielding, Neutron Shielding, Yb

1. INTRODUCTION

The group of rare earth elements (REs) is widely used in glass, ceramics, metallurgical industry, laser production, magnet production, oil catalyst and high-tech devices [1-6]. Ytterbium (Yb^{3+}), a rare earth (RE) element, has the potential to be utilized in the refinement of grains and the improvement of mechanical features of stainless steel products. With the half-life of 32 days, the radioactive isotope ^{169}Yb is sometimes employed as a radiation source in the area of nuclear medicine. Alloys of ytterbium-iron-cobalt-manganese are appropriate for the production of high-quality permanent magnets. Similarly, ytterbium-doped crystals have the ability to act as dopant materials in lasers.

Glass is a material of great interest to researchers due to its many benefits, including its perfect corrosion resistance, hardness, optical properties and ease of processing by various techniques. Rare earth doped glass matrices can be a valuable alternative to glass materials for potential use as scintillators and non-lead shields due to their low fabrication cost, chemical resistance, ease of fabrication and thermal stability properties [1-4]. Bismuth glass is in demand because of its extraordinary properties in amplifiers, lasers, optical data storage devices, etc. It is evident that the bismuth-borate matrix represents an ideal host matrix for REs, with applications spanning a range of disciplines including optical fibers, optical data storage devices, display technology, lasers, medical diagnostics and sensors [5-8]. Thus, Bi-doped glasses attract great attention from researchers [9-11]. Bi_2O_3 doped silica borotellurite glasses were studied by Geidam et al. [9] and it is suggested that the glasses are mechanically and thermally stable for applications in radiation shielding. Also, effect of Bi_2O_3 on gamma ray shielding and structural properties of borosilicate glasses recycled from high pressure sodium lamp glass was investigated and it was reported that good results were obtained [10]. Additionally, the Yb^{3+} ions present within the glasses serve as

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sensitizers for the various RE³⁺ ions, imparting the energy to increase the efficiency of emission, thereby imparting a dielectric nature to the glasses [8].

The use of glass doped with REs has become a key component in the development of photonic applications in modern life. In recent decades, there has been extensive investigation into the synthesis of materials doped with REs for application within lighting technologies. This is because of the necessity of improving the performance of solid-state devices, LEDs, lasers etc. [12]. It is also important to examine the influence of the external ionizing radiation on the glassy structures' microstructure. Learning about the interaction between radiation and glasses gives us enlightening information for the reaction of the glassy matrices to the radiation. The radiation features of the RE-doped glass matrix are strongly influenced by the composition of the host material and these features were commonly studied by the researchers [12-16]. In particular, it has been demonstrated in numerous studies that the addition of Yb³⁺ ions into glasses significantly enhances their capacity for effective protection [17-19]. The role of Yb₂O₃ in the radiation shielding properties of B₂O₃-TiO₂-BaO glasses was studied by Negm et al. [18], and it was obtained that this type of glass with low concentrations of Yb₂O₃ shows good shielding properties. Tamam et al. [19] reported that the photon shielding ability and dosimetry potential of the glass system improved with Yb³⁺ weight content of the glasses and they are recommended for gamma radiation inhibition functions such as nuclear waste container, structural and source shields.

In countries with developing economies, nuclear energy has become a significant aspect of the energy production landscape. The operation of the nuclear reactor will result in the release of highly penetrating and dangerous radiation. Accordingly, it is essential that the reactor site is safeguarded by an appropriate layer of protection. With the growth of the nuclear industry and the increasing utilisation of radiation technologies, it is therefore evident that the attenuation of radiation levels to ensure safety has become a crucial requirement. For this purpose, the goal of the study is to analyze the radiation-glasses interaction characteristics and to obtain the radiation attenuation parameters (RAP), which are significant to have the knowledge of attenuating abilities of the glasses, including the mass and linear attenuation coefficient values (MACVs and LACVs), mean free path value (MFPV), half value layer value (HVLV), electronic and atomic cross sections (ECS and ACS), effective and equivalent atomic number (Z_{eff} and Z_{eq}), effective electron density and conductivity (N_{eff} and C_{eff}), buildup factors (BUF), KERMA, fast neutron removal cross section value (FNRCVS), the mass stopping power (MSP) and projected range (R_p). In this context, the RAP of 50Bi₂O₃-15Li₂O-15PbO-(20-x)B₂O₃-Yb₂O₃ ($x = 0.5, 1.0, 1.5, 2.0$ and 2.5 mol%) glasses are evaluated by Phy-X/PSD [20], SRIM [21], PAGEX [22] and ESTAR (Institute) codes in different energies.

2. MATERIAL AND METHODS

The knowledge necessary for the computations of glass samples were taken from Ref. Bhemarajam et al. [8] to ascertain their potential for radiation shielding. The glasses with a composition of 50Bi₂O₃-15Li₂O-15PbO-(20-x)B₂O₃-Yb₂O₃ ($x = 0.5, 1.0, 1.5, 2.0$ and 2.5 mol%) have been obtained using the melt quenching technique. In melt quenching technique, firstly, the chemicals were weighted according to mol percent, then, they were ground to homogeneity using an agate mortar and placed into an alumina crucible. The obtained mixture is melted in an electric furnace at high temperatures. The desired melts were then quenched using a preheated brass plate and a quencher to produce a glass sample. Glass samples were annealed to avoid unwanted thermal stress and air bubbles in the synthesised glass [8]. Table 1 presents the densities, chemical compositions and molar volumes of the chosen glasses. The samples are labeled as Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5 for $x = 0.5, 1.0, 1.5, 2.0$ and 2.5 mol%, respectively. The study [8] outlines the methodology employed to prepare and assess the physical characteristics of the studied samples.

Table 1. The elemental components and density of the glasses [8].

Sample	Bi	O	B	Li	Pb	Yb	Density (g/cm ³)
Yb0.5	0.7295	0.1340	0.0147	0.0073	0.1085	0.0060	5.302
Yb1.0	0.7253	0.1333	0.0143	0.0072	0.1079	0.0120	5.403
Yb1.5	0.7213	0.1325	0.0138	0.0072	0.1073	0.0179	5.523
Yb2.0	0.7173	0.1318	0.0134	0.0071	0.1067	0.0238	5.635
Yb2.5	0.7133	0.1311	0.0129	0.0071	0.1061	0.0295	5.794

2.1. Radiation protection parameters for the interaction of charged and uncharged particles with matter

The Beer–Lambert statement can be written to determine MACVs,

$$I = I_0 e^{-\mu t} \tag{1}$$

$$\mu_m = \frac{\mu}{\rho} = \frac{\ln(I_0/I)}{\rho t} = \frac{\ln(I_0/I)}{t_m} \tag{2}$$

LACVs and MACVs are accepted as μ (cm⁻¹) and μ_m (cm²/g), respectively. The MACV is the potential for interaction between photons and the mass per unit area of a material, which provides insight into the absorption characteristics of the material in question. LACV is the fraction of incident photons attenuated per unit material thickness.

The HVLV refers to the capacity of radiation to penetrate a given substance. MFPV quantifies the average length traversed by radiation between two consecutive collisions within a given medium. HVLV and MFPV are calculated with the following equations

$$\text{HVLV} = \frac{\ln(2)}{\mu} , \quad \text{MFPV} = \frac{1}{\mu} \tag{3}$$

The interaction possibilities of electrons and atoms in a given volume of any given material is denoted by the symbols ECS and ACS, respectively. ACS (σ_a) and ECS (σ_e) can be acquired with the help of equations [23]

$$\text{ACS} = \sigma_a = \frac{N}{N_A} (\mu/\rho) , \quad \text{ECS} = \sigma_e = \frac{\sigma_a}{Z_{\text{eff}}} \tag{4}$$

A compound's atomic number is constituted by several elements; a compound with more than one element is referred to as an effective atomic number (Z_{eff}). Z_{eff} is obtained by [24]

$$Z_{\text{eff}} = \sigma_a / \sigma_e \tag{5}$$

The N_{eff} parameter is associated with the effective conductivity of the relevant material, which is dependent upon the excitatory photon energy. C_{eff} is in proportion to N_{eff} . N_{eff} and C_{eff} values can be written as

$$N_{\text{eff}} = \frac{\mu_m}{\sigma_e} , \quad C_{\text{eff}} = \left(\frac{N_{\text{eff}} \rho e^2 \tau}{m_e} \right) \times 10^3 \tag{6}$$

The N_{eff} value of the charged particles (CPs) can be determined by replacing MAC with the MSP value [21].

BUF, energy absorption build up factor (EABF) and exposure build up factor (EBF), can be found by the following equations below [25-27].

$$Z_{\text{eq}} = \frac{Z_1(\log R_2 - \log R) + Z_2(\log R - \log R_1)}{\log R_2 - \log R_1} \tag{7}$$

$$F = \frac{F_1(\log Z_2 - \log Z_{eq}) + F_2(\log Z_{eq} - \log Z_1)}{\log Z_2 - \log Z_1} \tag{8}$$

$$B(E, x) \begin{cases} 1 + \frac{(b-1)(K^x-1)}{(K-1)} & \text{for } K \neq 1 \\ 1 + (b-1)x & \text{for } K = 1 \end{cases} \tag{9}$$

$$K(E, x) = cx^a + d \frac{\tanh\left(\frac{x}{K_k}-2\right) - \tanh(-2)}{1 - \tanh(-2)} \quad \text{for } x \leq 40 \text{ mfp} \tag{10}$$

The R1 and R2 values represent the ratio of Compton scattering to total attenuation for adjacent elements with atomic numbers Z1 and Z2, measured in micrometres. F is a set of G-P FP (a, b, c, d, and XK coefficients) for the sample. F1 and F2 are the G-P FP for atomic numbers Z1 and Z2 at a specific energy. x and E refer to the penetration depth and primary photon energy, respectively.

In the case of glass systems, the KERMA with reference to air can be expressed as follows: [28]

$$\text{KERMA} = \frac{(\mu_{en}/\rho)_{\text{glass}}}{(\mu_{en}/\rho)_{\text{air}}} \tag{11}$$

The FNRCsv (ΣR) quantifies a material's potential capacity to attenuate neutrons and can be found by [20]

$$\Sigma R = \sum_i \rho_i (\Sigma R / \rho)_i \tag{12}$$

where ρ_i and $(\Sigma R / \rho)_i$ are the partial density of the glass and the mass removal cross-section of the *i*th constituent element, respectively.

SP can be formulated as [29]

$$\frac{dE}{dx} = 4\pi r_0^2 z^2 \frac{mc^2}{\beta^2} NZ \left[\ln \left(\frac{2mc^2}{I} \beta^2 \gamma^2 \right) - \beta^2 \right] \tag{13}$$

and for electrons the SP formula is

$$\frac{dE}{dx} = 4\pi r_0^2 \frac{mc^2}{\beta^2} NZ \left\{ \ln \left(\frac{\beta \gamma \sqrt{\gamma - 1}}{I} mc^2 \right) + \frac{1}{2\gamma^2} \left[\frac{(\gamma - 1)^2}{8} + 1 - (\gamma^2 + 2\gamma - 1) \ln 2 \right] \right\} \tag{14}$$

The Rp value is assumed as

$$\text{Rp}(E) = \int_0^E \frac{1}{\left(-\frac{dE}{dx} \right)} dE \tag{15}$$

3. RESULTS and DISCUSSION

3.1. Radiation shielding analysis of charged and uncharged particles

The MACVs varied with photon energies were illustrated in Fig. 1(a). The MACRs exhibited a sharp reduction due to the photoelectric (PE) process at low energies (1-100 keV). At mid-energies (100 keV-5 MeV), they exhibited a slight change due to Compton scattering (CS). With increasing energy, the MACRs exhibited an increase due to pair production (PP) at high energies (>5 MeV) [30]. It can be said that the glasses protect well and Yb2.5 is slightly better than other results. The glasses shield equally well. Additionally, the LACVs that vary with photon energies were presented in Fig. 1(b), and the same shielding feature was found for the glasses. In Fig. 2, it is demonstrated that the MACVs of the glasses have better shielding features especially at low energies than the previously investigated shield materials [31].

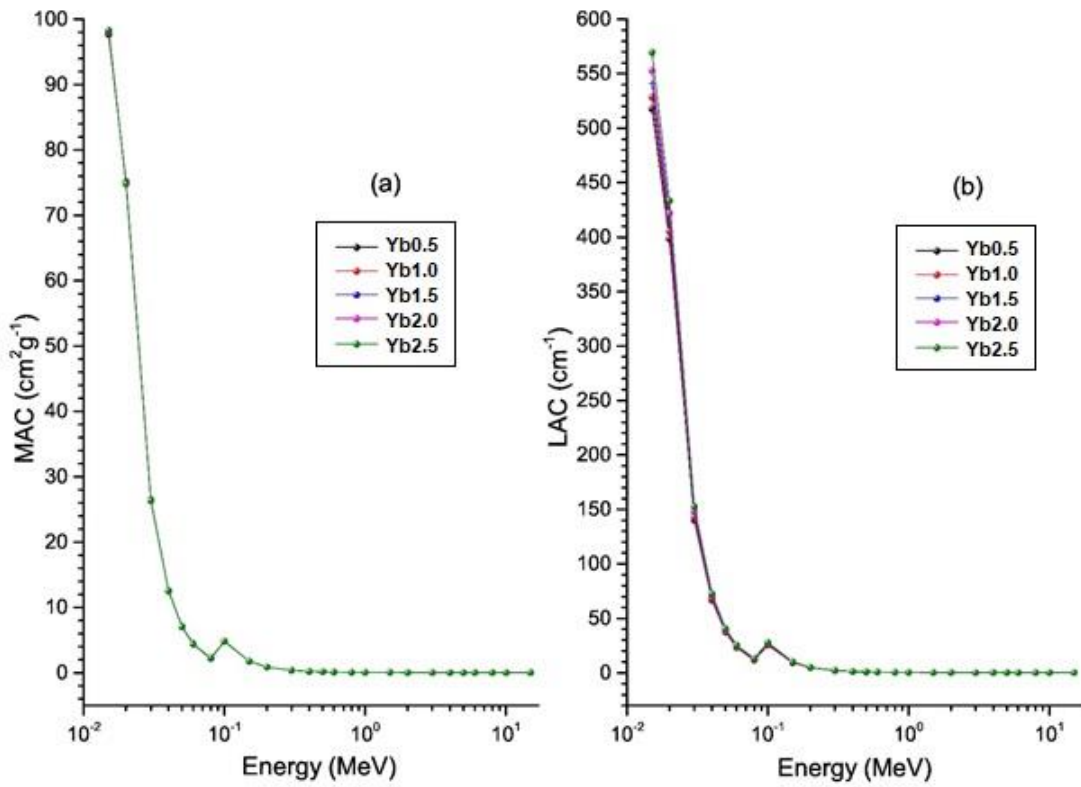


Figure 1. Variations of MAC (a) and LAC (b) values with photon energies.

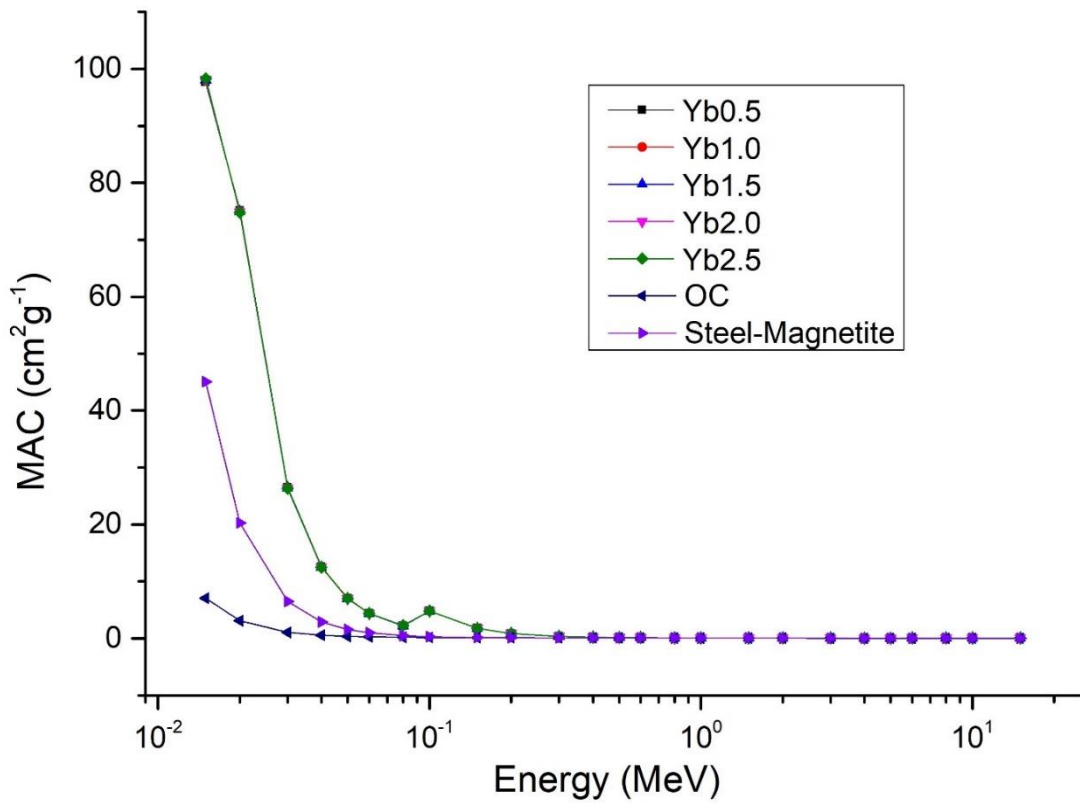


Figure 2. Comparison of MAC values of the glasses with previously studied shields.

The MFP and HVL results were exhibited in Figs. 3(a)-(b). At energies dominated by CS, photons were observed to exhibit a high probability of scattering. Consequently, thicker materials were found to be necessary, as they exhibited a lower probability of absorption and a greater photon MFPV. In the high energy region, lower HVL and MFP results were observed to be beneficial for enhanced shielding capability. In the high-energy region, lower HVL and MFP values are required, while in the lower energy region, thicker materials are necessary. In this regard, the MFP and HVL results are ordered as follows: $Yb2.5 < Yb2.0 < Yb1.5 < Yb1.0 < Yb0.5$. It is worth noting that glass Yb0.5 has the lowest level of shielding ability due to its high HVLVs and MFPVs, and Yb2.5 may offer the highest level of shielding among the glasses due to the lowest values. As seen in Fig. 4, the HVLVs of the studied glasses are lower than those of the materials reported before [31], so it can be said that the glasses are better protection materials.

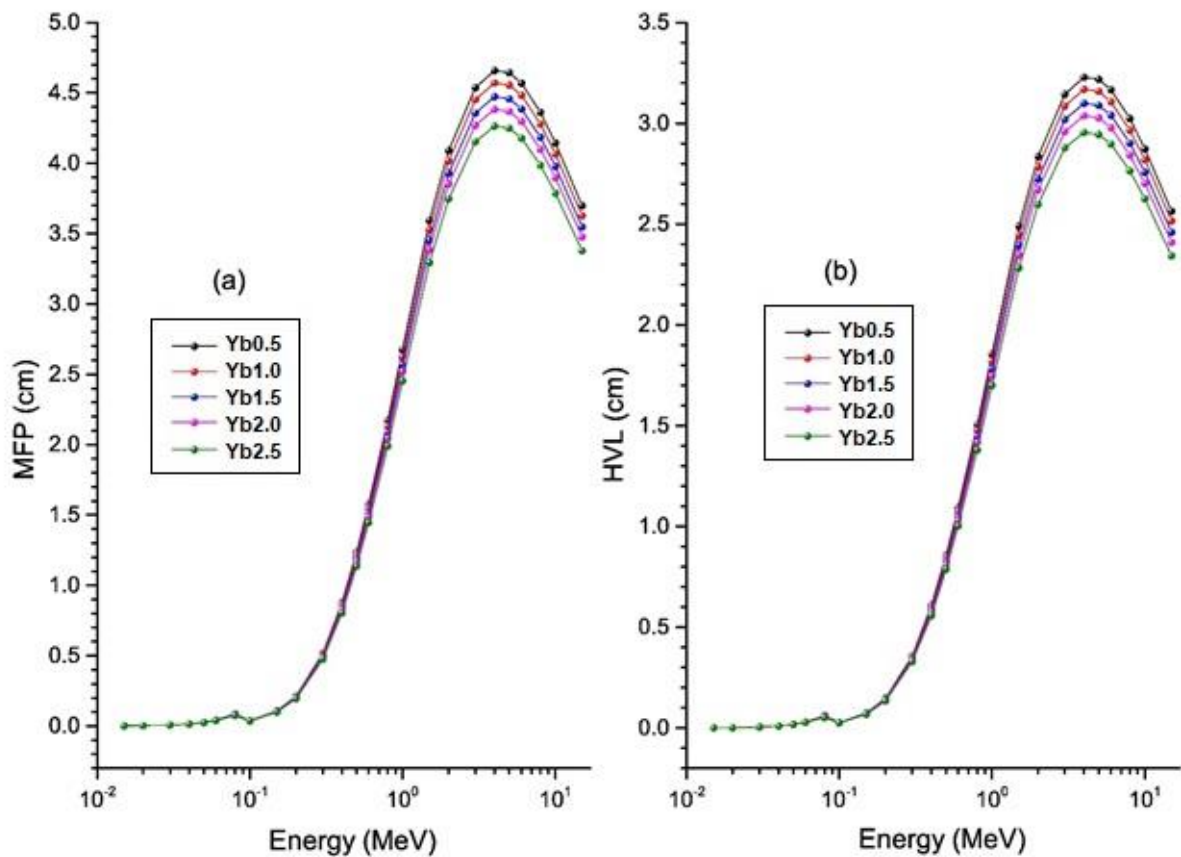


Figure 3. Variations of MFP (a) and HVL (b) values with photon energies.

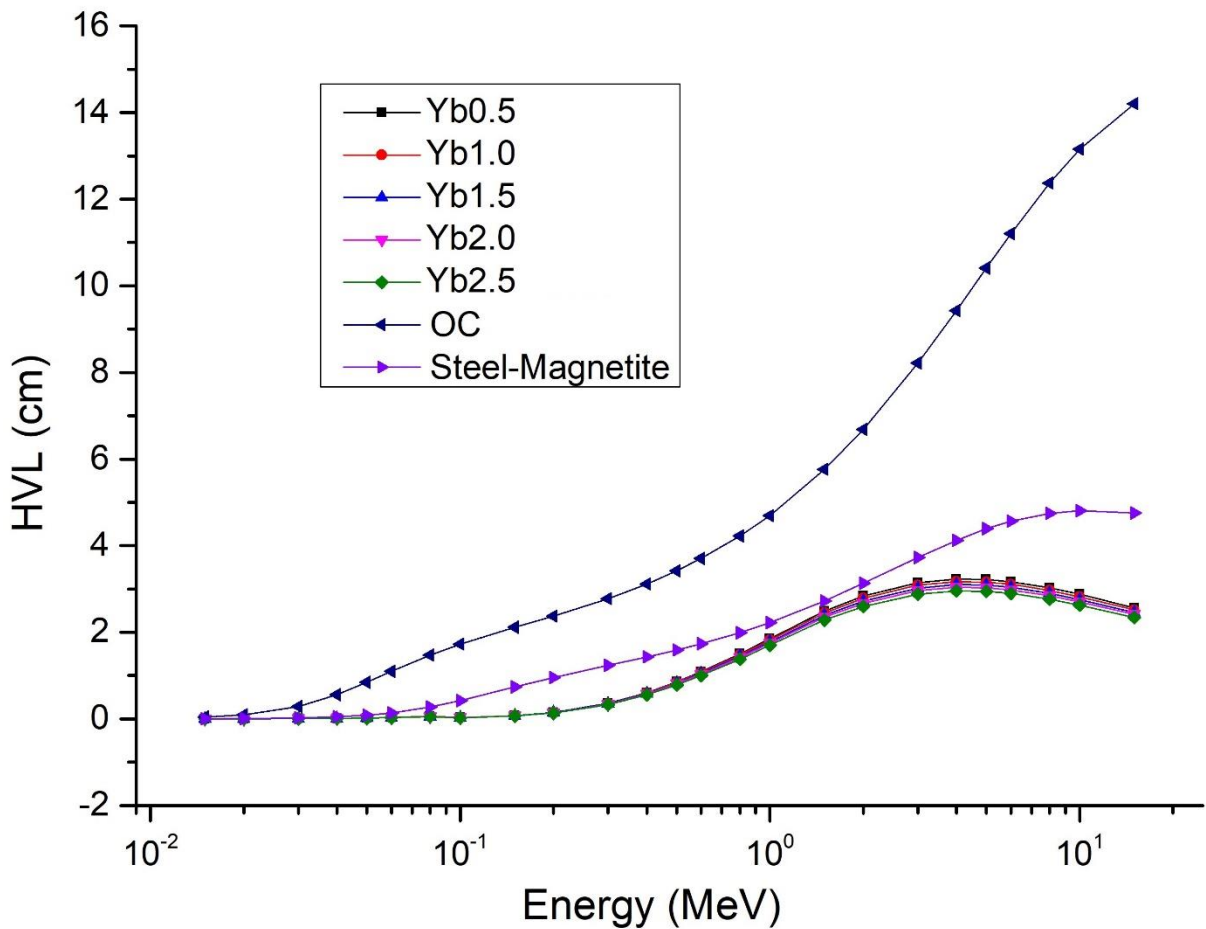


Figure 4. Comparison of HVLs of the glasses with previously studied shields.

It can be of interest to enquire further into the interaction probabilities of a material, which are shown by ACS and ECS in Figs. 5(a) and (b). The results demonstrate that the ACS and ECS values exhibited a correlation with the energy of the photon. It can be postulated that the higher ACS and ECS results point to a better protective glass. It would appear that Yb2.5 has the highest protection character among all the glasses, based on the ACS and ECS values.

The Z_{eff} results, which is displayed in Fig. 6(a), gave the highest values because of the PE at low energy levels. The results exhibited a sharp decline and subsequent increase as the energy increases, before stabilizing at high energies. Glass with Yb2.5 composition, which contains a higher amount of Yb, shows the highest Z_{eff} values and thus the highest shielding performance. On the other hand, Yb0.5 with a lower Yb content show the lowest Z_{eff} values and the lowest protection ability. The Z_{eq} values of the glasses are displayed in Fig. 6(b). Yb2.5 presents a greater interaction, while Yb0.5 shows the least interaction between radiation and matter. The obtained values indicate that Yb2.5 exhibits a more pronounced interaction with radiation than Yb0.5, which demonstrates the least interaction among the glasses investigated. The Z_{eff} values of the glasses are also compared with those of the previous results reported for widely used shielding materials by Bashter [31]. The values are higher than the ordinary concrete (OC) and steel-magnetite, thus the protection abilities are higher.

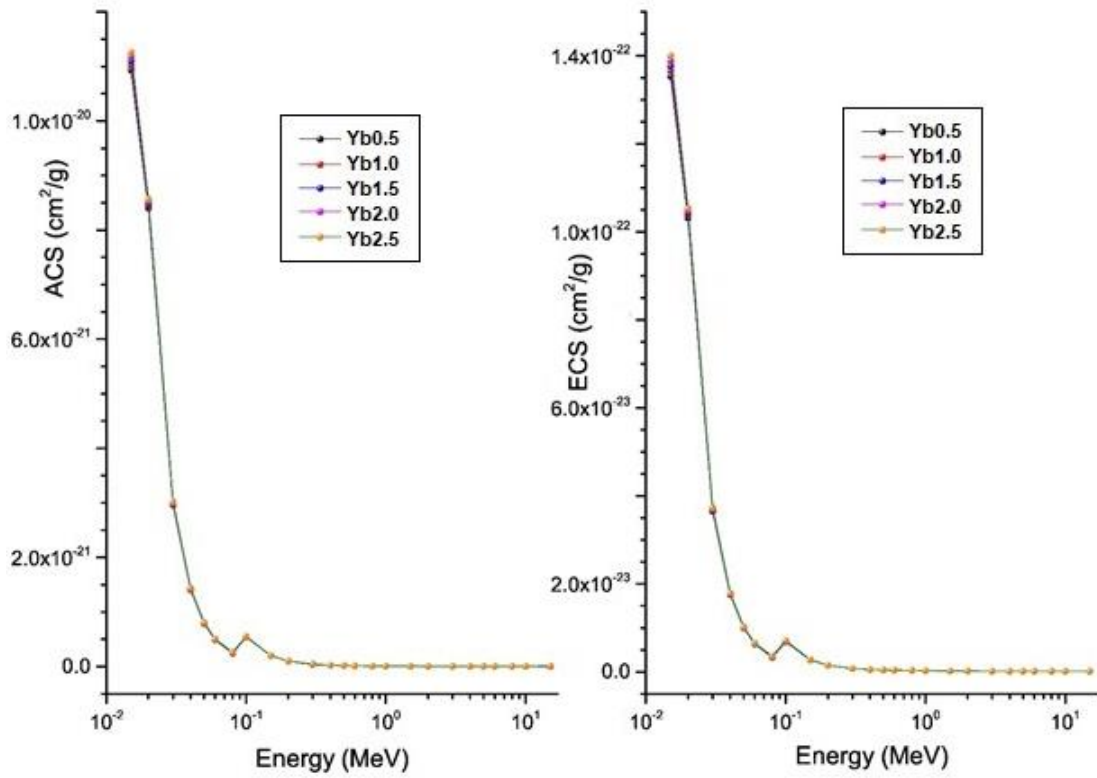


Figure 5. Variations of ACS (a) and ECS (b) values with photon energies.

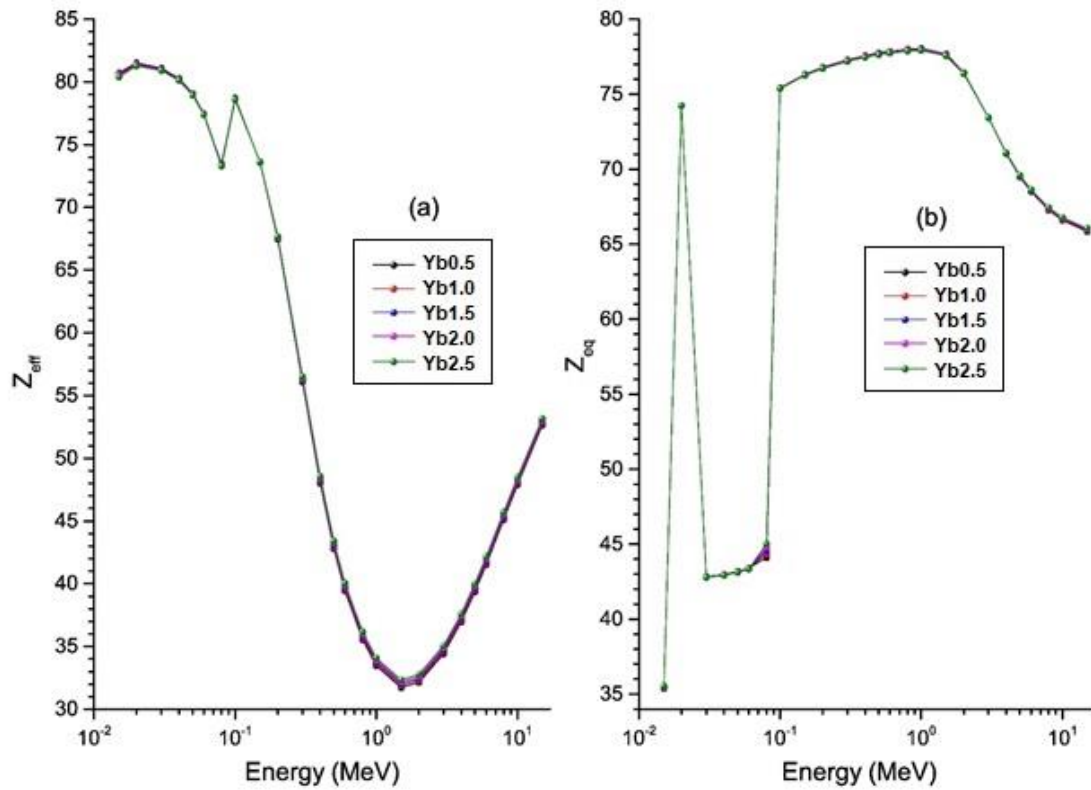


Figure 6. Changes of Z_{eff} (a) and Z_{eq} (b) values with photon energies.

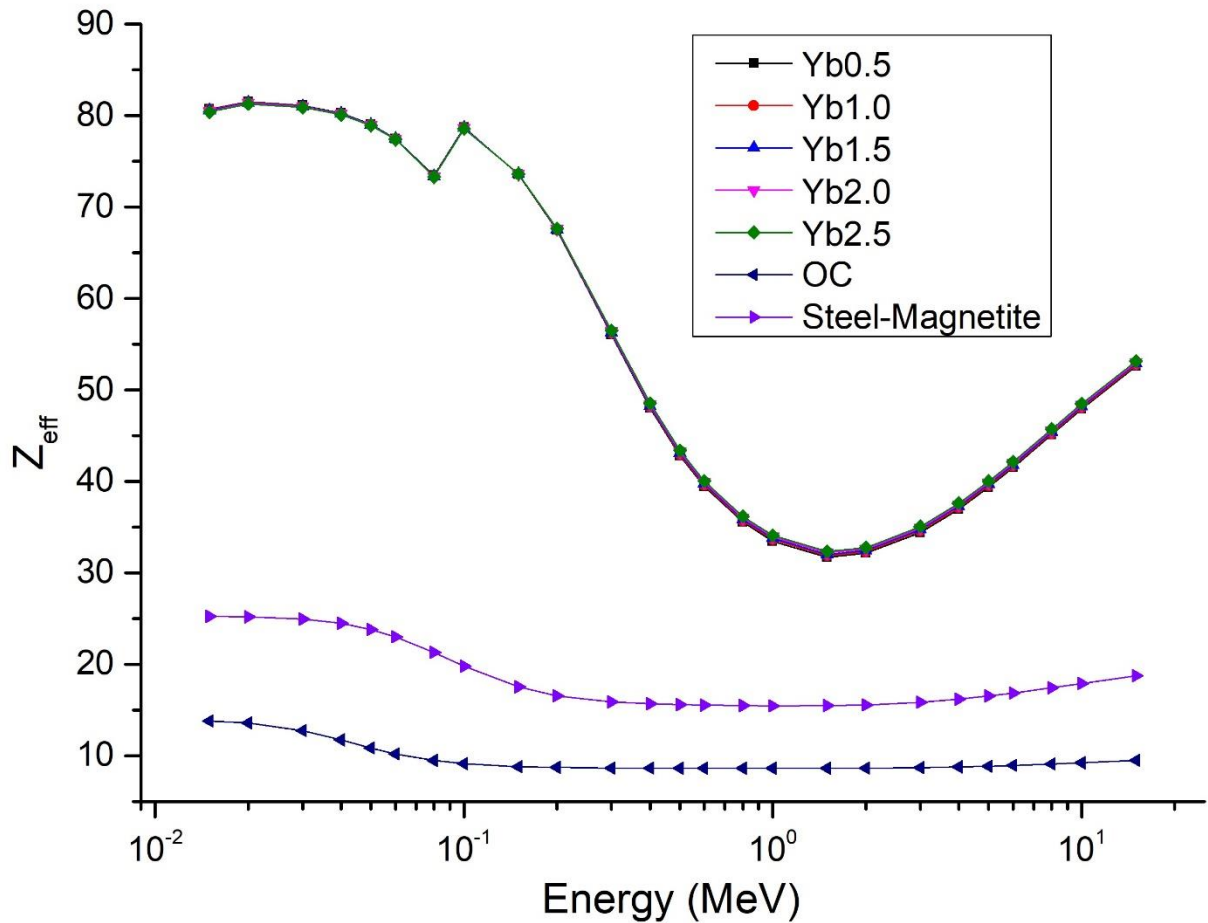


Figure 7. Comparison of Z_{eff} values of the glasses with previously studied shields.

Photon-material interactions result in alterations to the number of free electrons within the material, influenced by PE, CS, and PP interactions. This change is contingent upon the N_{eff} results, which correspond to the number of conduction electrons within the material. Consequently, C_{eff} is proportional to the N_{eff} of the glasses. The alterations in C_{eff} and N_{eff} results according to photon energies are illustrated in Figs. 8(a)-(b). It is noticed that Yb2.5 exhibits the highest C_{eff} and N_{eff} results among the glasses. The very sharp peak seen for Z_{eff} , C_{eff} and N_{eff} around 0.1 MeV can be attributed to the K absorption edges of the Bi and Pb elements [32].

The EABF and EBF of the glasses for 16 penetration depths were defined by the help of Phy-X/PSD and their changes with the incident energy are shown in Figs. 9 and 10, respectively. In the PE region, all energies absorb lower energy photons, resulting in small buildup factors. A large number of scattered photons in the CS region can cause the photon accumulation to increase, causing the factors to reach the highest values in the mid-energy region. The photon absorption rate is high in the PP region, which results in lower buildup factors in higher energy regions. The results indicate that the photon cluster for Yb0.5 is higher than that of the other glasses, as evidenced by the values obtained for EBF and EABF. Additionally, sharp increases in the energies at ≈ 0.03 and 0.1 MeV can be explained as L1 and K absorption edges of Bi [32].

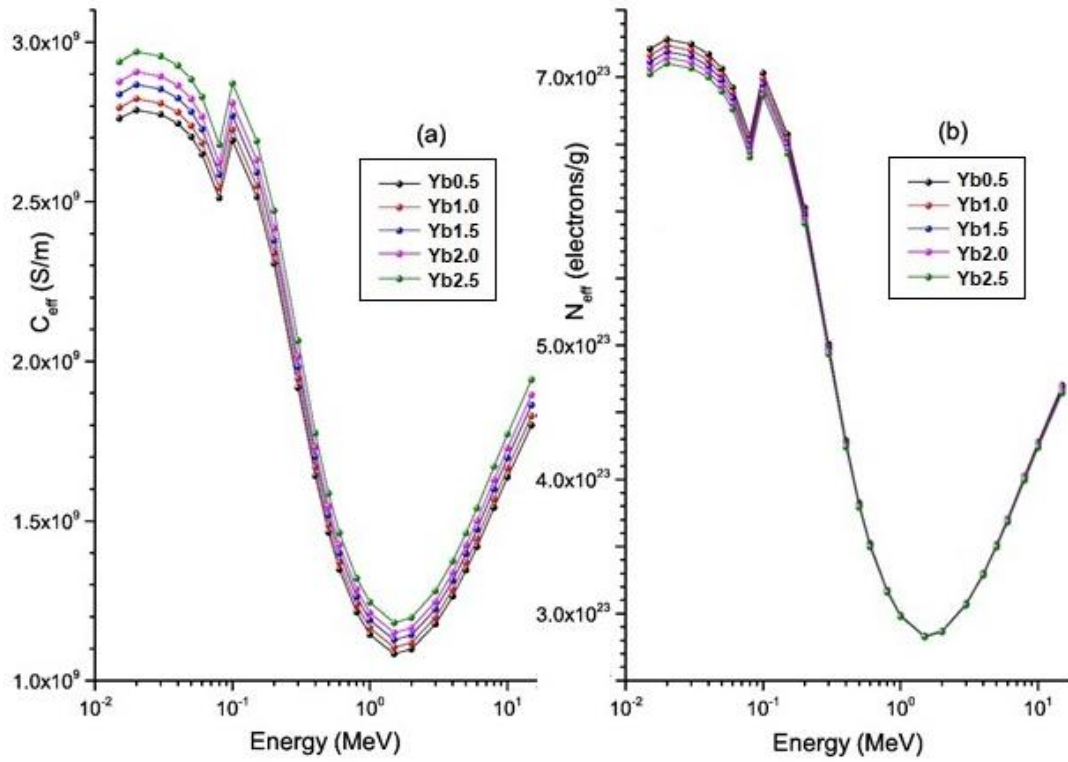


Figure 8. Changes of C_{eff} (a) and N_{eff} (b) values with photon energies.

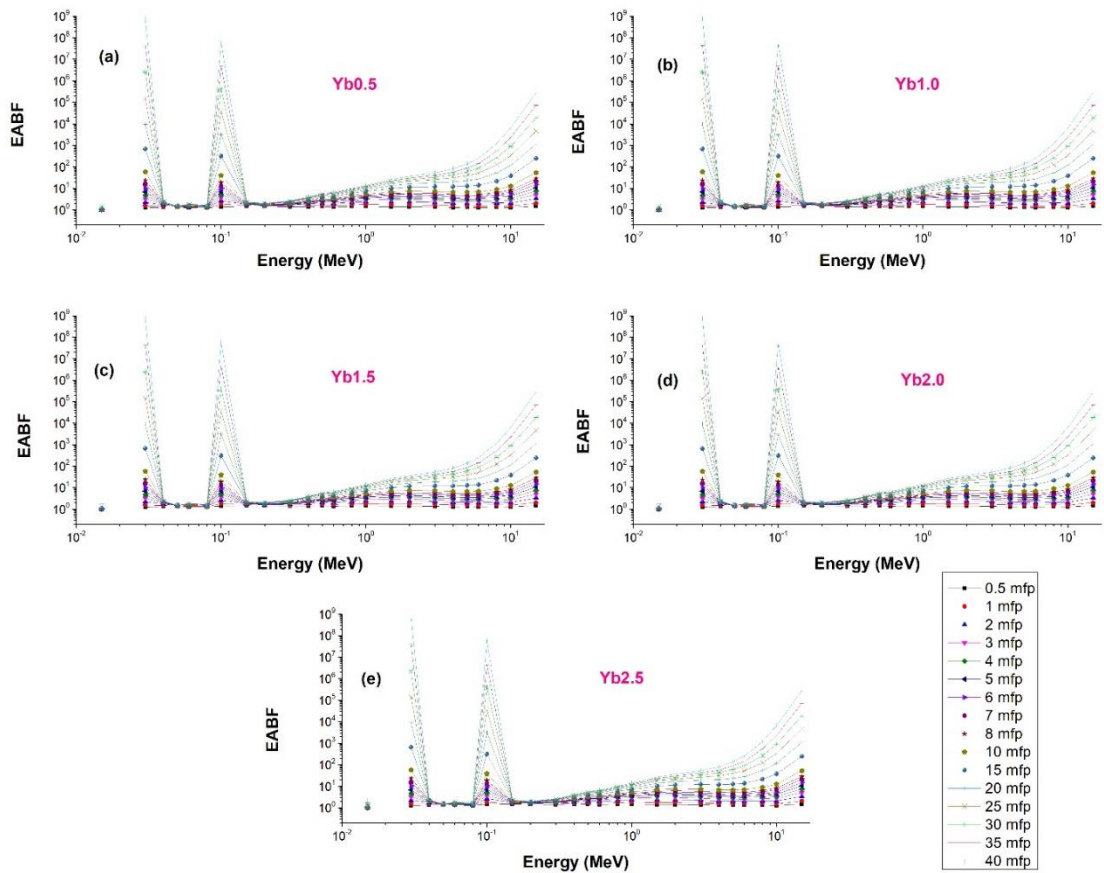


Figure 9. Variations of EABF values with photon energies.

The μ_{en} values are estimated and demonstrated in Fig. 11(a). The results (in cm²/g) are in the range 0.022946 - 2266.293, 0.022941 - 2261.435, 0.022936 - 2256.637, 0.022932 - 2251.895 and 0.022927 - 2247.195 for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. The findings follow this order: Yb2.5 < Yb2.0 < Yb1.5 < Yb1.0 < Yb0.5. Additionally, the KERMA values are determined and shown in Fig. 11(b). At low energies, KERMA values are at their minimum, and begins to increase and reaches its highest value at 0.04 MeV. The data then demonstrated a rapid decline, followed by an average behavior. Yb0.5 achieved the greatest KERMA at 0.04 MeV while Yb2.5 had the lowest KERMA.

The neutron attenuation of the glasses is determined by means of the thermal and FNRCS values and shown in Fig. 12. Yb2.5 has the highest value and Yb0.5 the lowest for fast neutrons. The FNRCS values are compared with some shielding materials such as Bi₂O₃ glass, graphite and borogypsum reported previously and illustrated in Fig. 12 [30,33,34]. Unlike fast neutrons, thermal ones have the lowest value for Yb2.5 and the highest for Yb0.5. This phenomenon can be explained by the high content of heavy elements, for example bismuth, which do not have the capacity to absorb fast neutrons effectively within the glass composition. Given that B is an effective thermal neutron absorber [35], the elevated B₂O₃ concentration in the glass Yb0.5 has resulted in a notable enhancement in the glass's thermal neutron absorption capacity. It can be also seen that thermal neutron attenuation is better for Yb⁺³ doped glasses than that of cement mortar with 10% and 20%, replacement of borosilicate glass powder studied by Jang et al. [36].

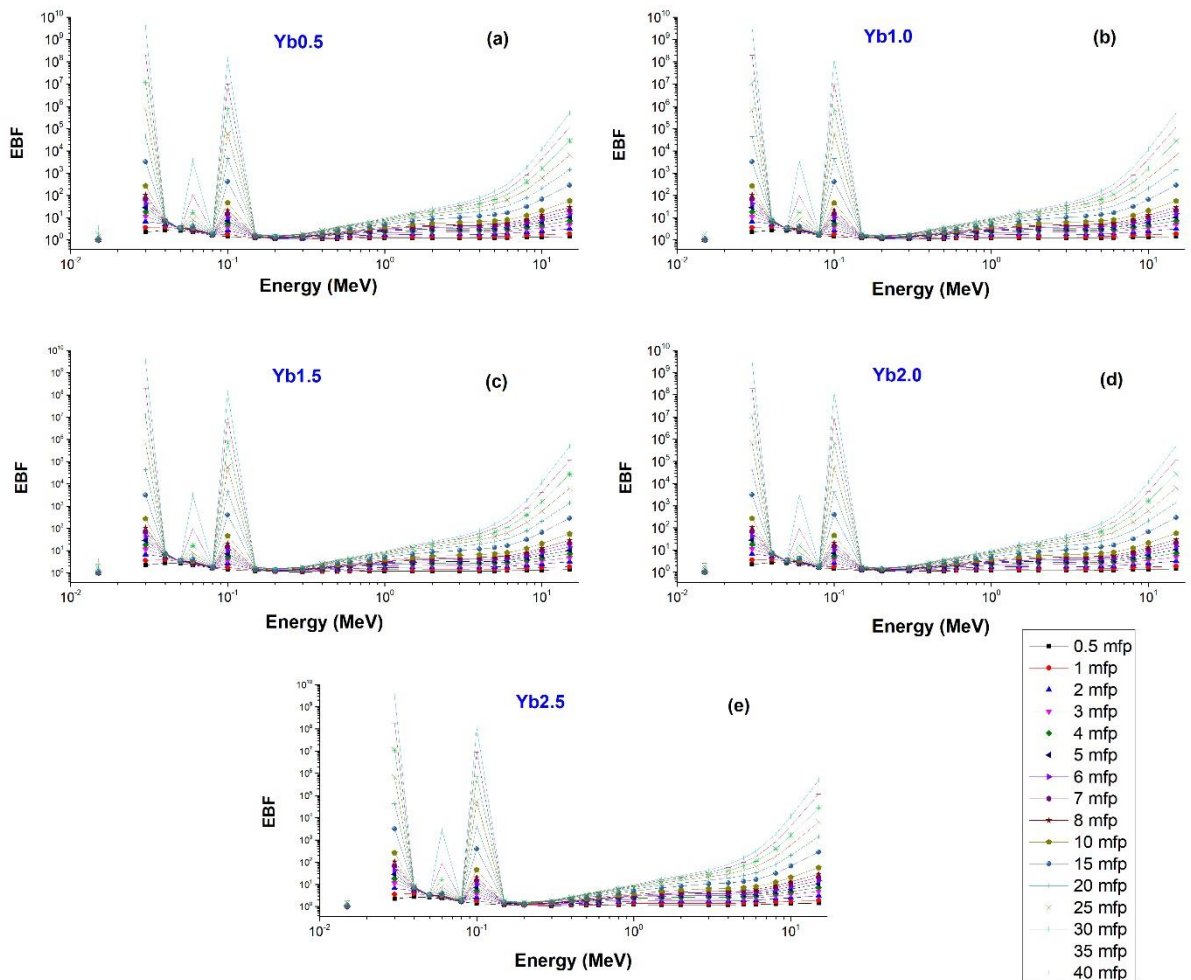


Figure 10. Variations of EBF values with photon energies.

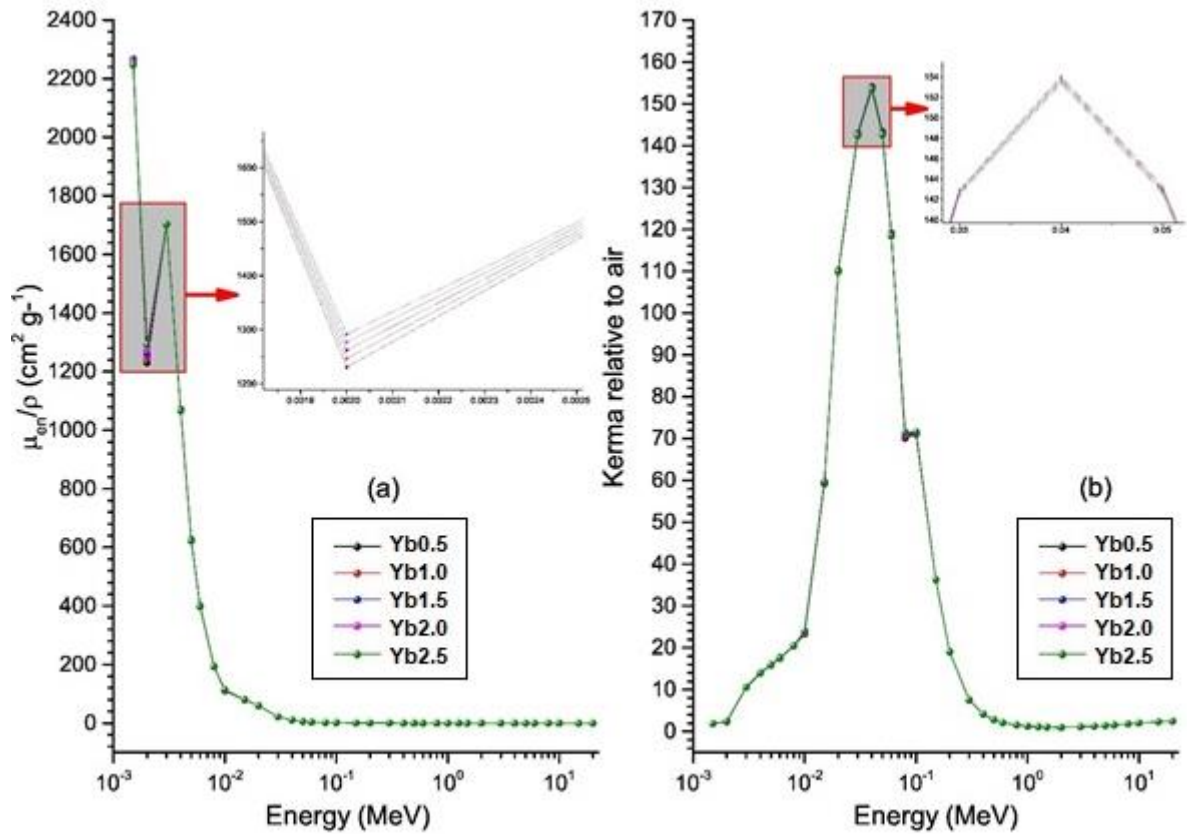


Figure 11. Variations of μ_{en}/ρ (a) and KERMA relative to air (b) values with photon energies.

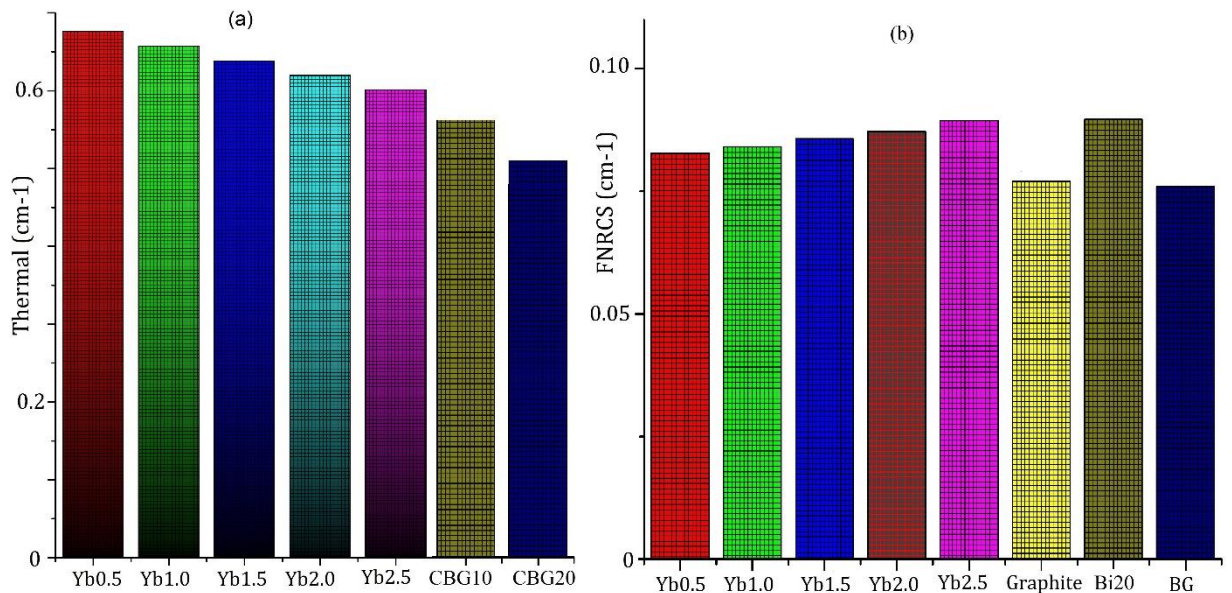


Figure 12. Comparison of thermal and FNRCS values of the glasses as comparative with previously reported materials.

The interaction of CPs with glasses is of paramount importance in radiation therapy. In order to ascertain the suitability of glasses for this aim, the MSP values of the electron, alpha and proton have been estimated for the glasses in question and the results have been presented in Fig. 13. The MSP values of electrons exhibit a decrease with increasing energy, followed by a gradual increase. In this case, the energy

loss mechanism of electrons is taken into account. It is supposed that electrons experience Coulomb interactions and bremsstrahlung. The MSP values (in MeV cm²/g) of electrons are among 1.196 – 137.0, 1.195 – 137.1, 1.194 – 137.2, 1.193 – 137.3 and 1.192 – 137.3 for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. MSP values for protons, increase with energy and subsequently, the value reduces in a gradual manner until the maximum is reached. The MSP results (in MeV cm²/g) of proton are among 1.306 – 219.136, 1.306 – 218.436, 1.305 – 217.735, 1.304 – 217.034 and 1.304 – 216.334 for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. The MSP for alpha particle exhibits an increase in energy prior to a subsequent decrease. It should be noted that the behavior of proton and alpha particles can be affected by the dependence of their electronic and nuclear energy losses, that in turn, can be influenced by the kinetic energy of the particles in question. The MSP values (in MeV cm²/g) of alpha are among 8.881 – 620.172, 8.877 – 618.671, 8.872 – 617.169, 8.868 – 615.668 and 8.863 – 614.166 for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. The MSP values for the glasses are given in the following order for the CPs: MSP_{electrons} < MSP_{protons} < MSP_{alpha}.

The Rp results are calculated by performing the SRIM and the ESTAR codes at 0.01 - 30 MeV, and are displayed in Fig. 14. The Rp values (in μm) for electron are among 1.1976 – 22180, 1.1765 – 21765, 1.1522 – 21310, 1.1306 – 20887 and 1.1007 – 20331 for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. The Rp results (in μm) for proton are among 0.1082 – 3380, 0.1063 – 3320, 0.1042 – 3250, 0.1022 – 3180 and 0.0995 – 3100 for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. The Rp results (in μm) for alpha are among 0.0713 – 157, 0.0701 – 152, 0.0686 – 148, 0.0673 – 144 and 0.0655 – 125 for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. The order of Rp values for the glasses is as follows: R_{alpha} < R_{proton} < R_{electron}. Additionally, the largest Rp values of CPs follow the following order: Yb0.5 > Yb1.0 > Yb1.5 > Yb2.0 > Yb2.5. It is of significant importance to note that the glass with the highest density exhibits a lower Rp value, while the glass with the lowest density exhibits the greatest Rp value.

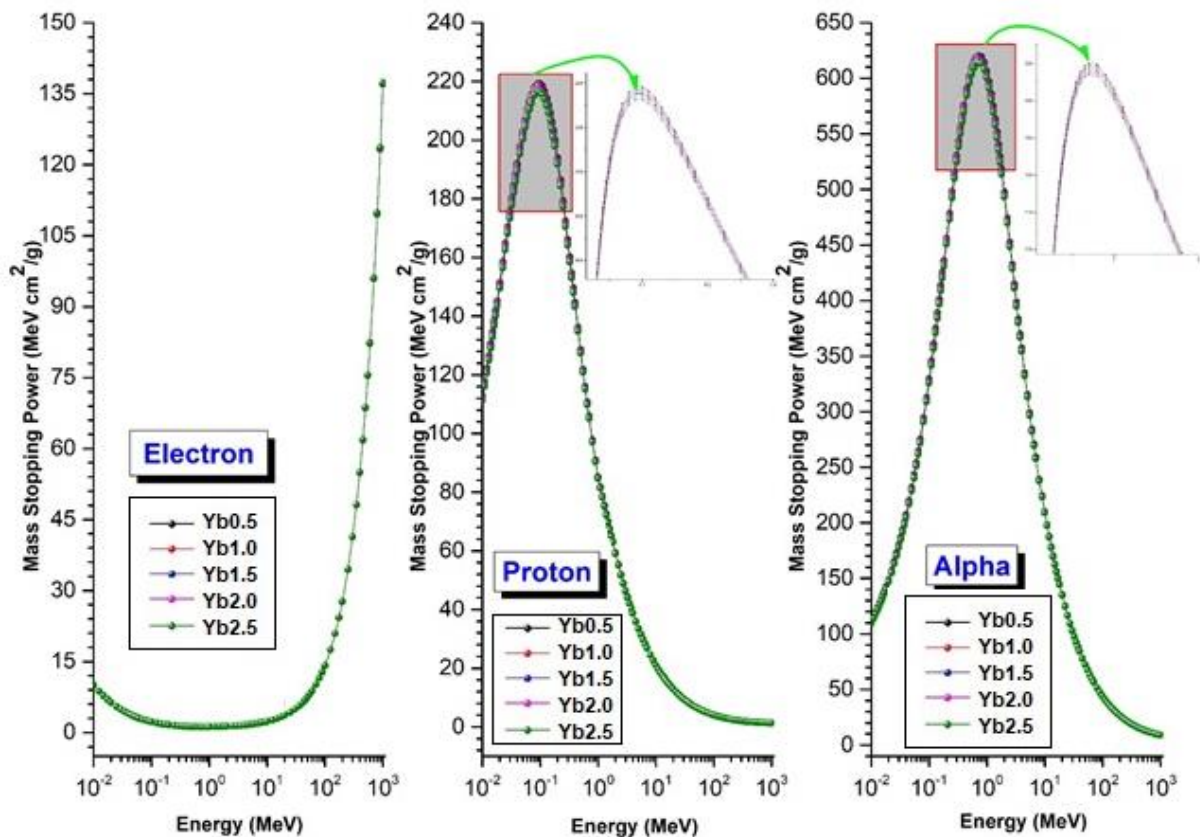


Figure 13. Dependence of MSP values of alpha, proton and electron particles versus the energy.

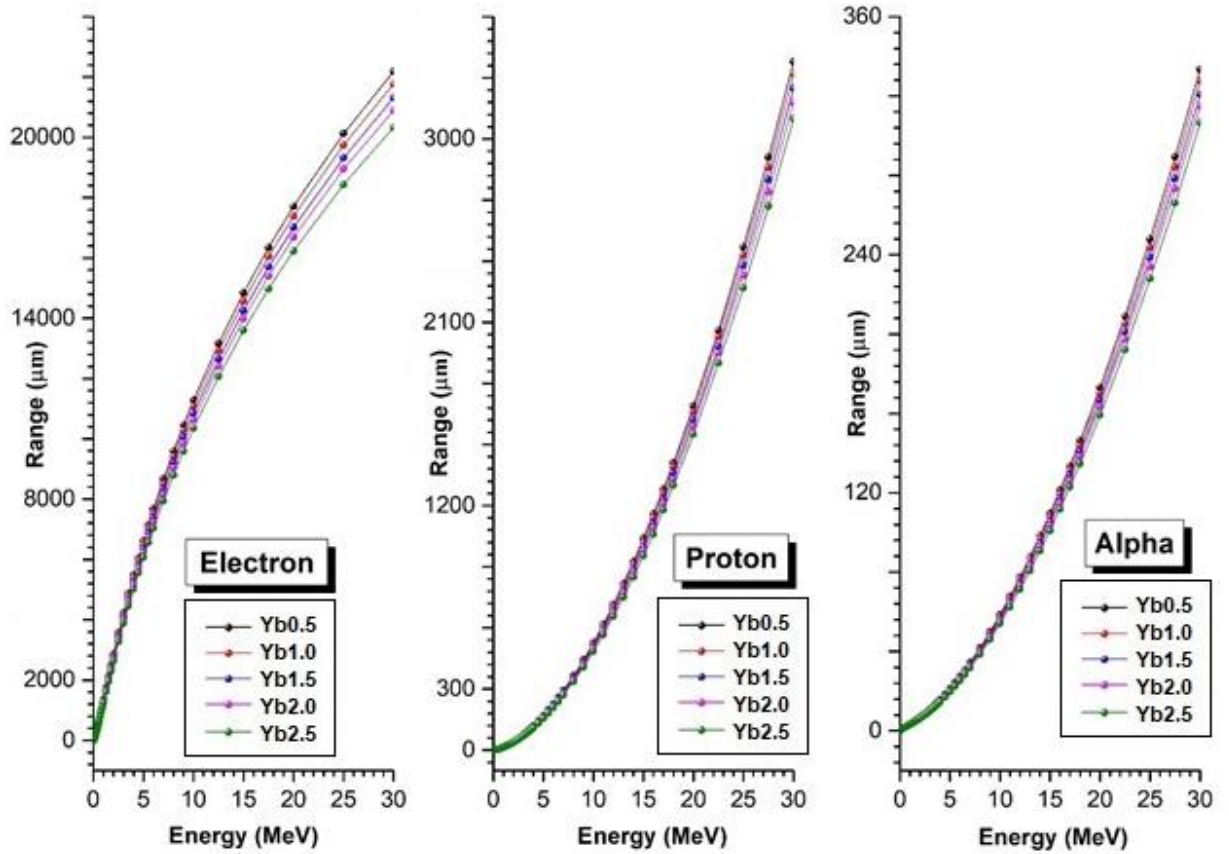


Figure 14. Changes of R_p values of alpha, proton and electron particles versus energy.

The Z_{eff} results of the glasses are illustrated in Fig. 15. The Z_{eff} values for electron are among 21.422 - 40.659, 21.623 - 40.804, 21.778 - 40.948, 21.925 - 41.091, 22.069 - 41.217 for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. The Z_{eff} values for proton are among 12.185 - 24.951, 12.191 - 25.814, 12.196 - 26.041, 12.202 - 26.164, 12.207 - 26.265 for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. The Z_{eff} results for alpha are among 13.447 - 38.556, 13.469 - 38.566, 13.703 - 38.576, 13.726 - 38.586, and 13.749 - 38.596 for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. From the results, the largest Z_{eff} values of the glasses for all the CPs give the following order: Yb0.5 < Yb1.0 < Yb1.5 < Yb2.0 < Yb2.5. The Z_{eff} values for electrons increase gradually with the kinetic energy for each sample. It has been documented that there is a non-monotonic variation in the behavior of alpha particles and protons in both low- and high-energy contexts. Furthermore, Z_{eff} values tend to increase with increasing glass density.

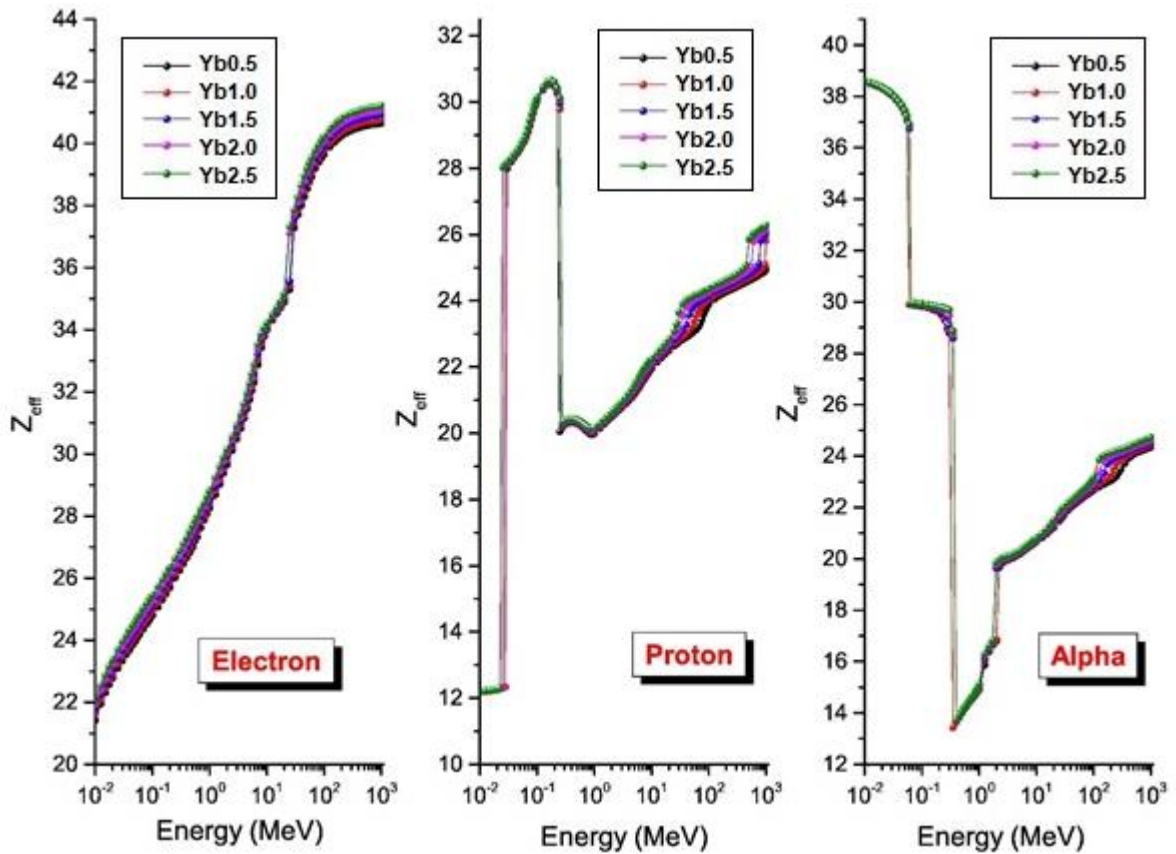


Figure 15. Variations of Z_{eff} values for alpha, proton and electron interactions with the studied glasses.

Also, the N_{eff} results of the glasses for the CPs are calculated and demonstrated in Fig. 16. The N_{eff} results (electrons/g) of electron are among $1.913 \times 10^{23} - 3.632 \times 10^{23}$, $1.921 \times 10^{23} - 3.624 \times 10^{23}$, $1.923 \times 10^{23} - 3.617 \times 10^{23}$, $1.926 \times 10^{23} - 3.609 \times 10^{23}$, $1.927 \times 10^{23} - 3.600 \times 10^{23}$ for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. The N_{eff} values (electrons/g) of proton are among $1.088 \times 10^{23} - 2.729 \times 10^{23}$, $1.083 \times 10^{23} - 2.716 \times 10^{23}$, $1.077 \times 10^{23} - 2.703 \times 10^{23}$, $1.071 \times 10^{23} - 2.691 \times 10^{23}$, $1.066 \times 10^{23} - 2.678 \times 10^{23}$ for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. The N_{eff} values (electrons/g) of alpha are among $1.201 \times 10^{23} - 3.444 \times 10^{23}$, $1.196 \times 10^{23} - 3.426 \times 10^{23}$, $1.210 \times 10^{23} - 3.407 \times 10^{23}$, $1.205 \times 10^{23} - 3.389 \times 10^{23}$, $1.201 \times 10^{23} - 3.371 \times 10^{23}$ for Yb0.5, Yb1.0, Yb1.5, Yb2.0 and Yb2.5, respectively. A similar pattern is observed in the N_{eff} values as Z_{eff} for electrons, which increase gradually with increasing kinetic energy for each glass sample. It is observed that non-monotonic variations in N_{eff} values for alpha and proton are present in both low and high-energy regions.

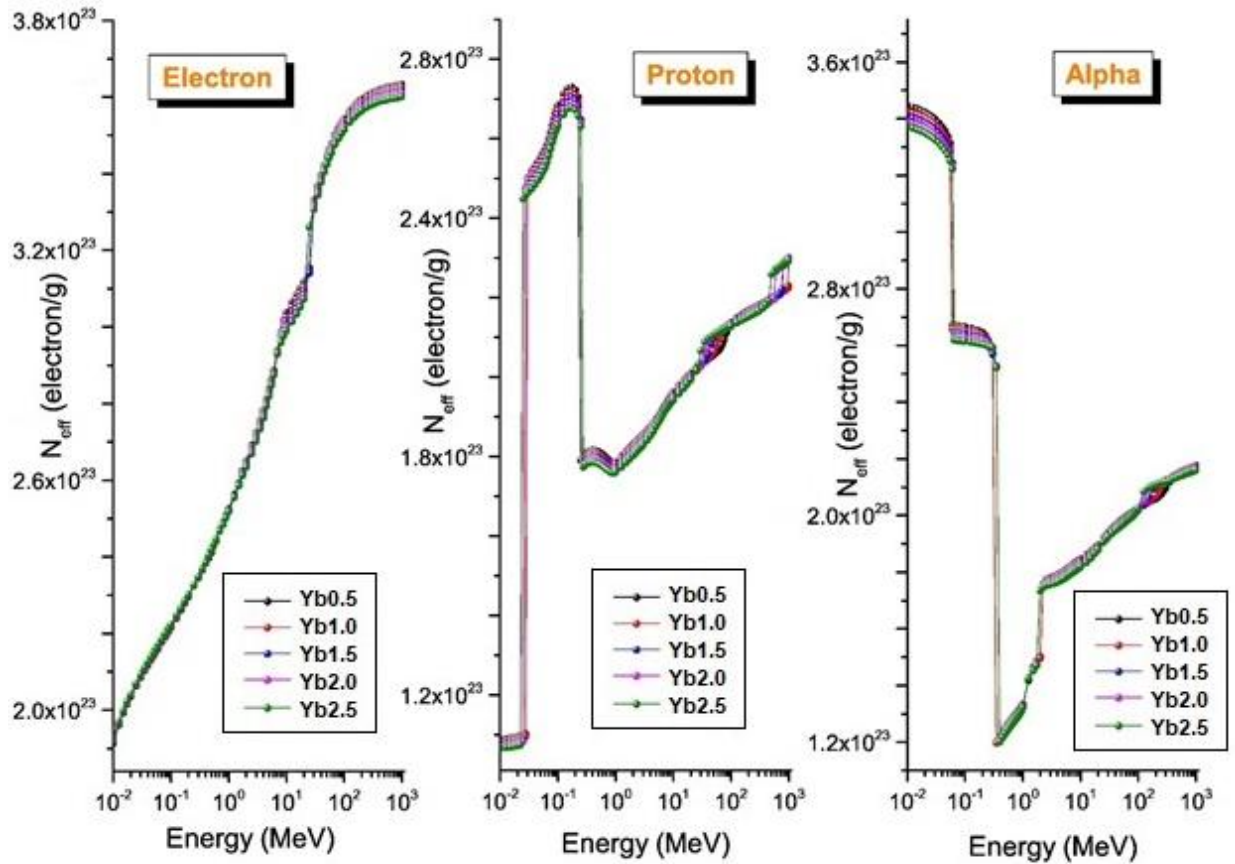


Figure 16. Variations of N_{eff} values for alpha, proton and electron interactions with the studied glasses.

4. CONCLUSIONS

In the paper, the radiation protection properties of Yb³⁺ doped glasses with contents of 50Bi₂O₃-15Li₂O-15PbO-(20-x)B₂O₃-Yb₂O₃ (x = 0.5, 1.0, 1.5, 2.0 and 2.5 mol%) were examined for fast neutron, CPs and gamma-ray by Phy-X/PSD, PAGEX, ESTAR and SRIM codes. The MACVs, LACVs, MFP, HVL, ACS, ECS and Z_{eff} values were evaluated and it was observed that these parameters exhibited energy-dependent behavior due to the photon interaction processes occurring across a broad spectrum of energies. In addition, the highest and lowest MACVs, LACVs, ACS, ECS, Z_{eff} and μ_{en} values were found for Yb_{2.5} and Yb_{0.5} while those of MFPVs and HVLVs were determined for Yb_{0.5} and Yb_{2.5} glasses, respectively. Since the best neutron shielding value was obtained for Yb_{2.5} that has the highest FNRCs value among the glass samples, it can be explained that it is the glass with the best neutron shielding feature. The CPs' MSP results for the glasses are obtained in the order of alpha>proton>electron. The glass with the Yb_{0.5} composition recorded the longest R_p values for CPs, while the glass with Yb_{2.5} gave the shortest values. This difference can be assigned to the necessity of minimum thickness. The glasses subjected to investigation displayed an enhanced capability to shield against a range of CPs when Yb³⁺ was introduced. The study revealed that Yb_{2.5} offers the highest level of protection against gamma rays, neutrons and CPs, while Yb_{0.5} demonstrated the least effective performance among the glasses. It can be concluded that glasses with a higher proportion of Yb³⁺ provide a superior level of protection, while those with a lower proportion of Yb³⁺ offer a lesser degree of protection. As a result of this, it can be said that Yb³⁺ activated glasses with a composition of 50Bi₂O₃-15Li₂O-15PbO-(20-x)B₂O₃-Yb₂O₃ can be applied as good shields for use in a wide range of applications. The results suggest that the investigated glasses with high Bi (Bi50) content and Yb³⁺ addition may have excellent radiation attenuation properties making them a promising option for radiation shielding applications. The attenuation of the RE-doped glasses is

a notable factor in achieving enhanced shielding, which in turn makes them appropriate for use in lighting technologies. It can thus be posited that Yb³⁺-doped glasses will prove instrumental in the advancement of material designs.

Declaration of Ethical Standards

Not applicable

Credit Authorship Contribution Statement

M. AYGUN: Conceptualization, Methodology, Writing original draft, Visualization, Investigation.

Z. AYGUN: Methodology, Writing original draft, Visualization, Investigation.

I. HAN: Writing original draft, Visualization, Investigation.

E. NARMANLI HAN: Writing- review, Visualization and editing.

Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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Not applicable

Data Availability

Data of this study are available from the corresponding author on responsible request.

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