

Seiberg-Witten-Like Equations on 8-Manifolds without Self-Duality

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Abstract

In this paper, Seiberg–Witten–like equations without self–duality are defined on 8 –dimensional manifolds. Then, non–trivial and flat solutions are given to them on \mathbb{R}^8 . Finally, on 8 –real–dimensional Kähler manifolds a global solution to these equation is obtained for a given negative and constant scalar curvature.

Keywords: Seiberg–Witten equations; *Spin* and *Spin^c* geometry; Curvature; Without self–duality.

1. Introduction

On the 4 –manifolds, Seiberg–Witten equations introduced by E. Witten are consisted of Dirac equation and Curvature equation [9]. These equations provide information about the topology and geometry of the 4 –manifolds [3,4,6,8,9]. To define these equations, one needs two entities as an $i\mathbb{R}$ valued connection 1 –form and a spinor field. Dirac equation can be defined on any manifold endowed with *Spin^c* –structure. But, defining the curvature equation needs self–duality concept of two form. Since self–duality concept is meaningful only in 4 –dimension, generalized self–duality concept is given to define the curvature equation on a four– dimensional manifold. Accordingly, Seiberg–Witten equations are investigated up to 4 –dimensional manifolds by defining generalized self–duality concept [1,5]. On 8 –manifolds, Seiberg–Witten–like equations have been studied in [1,2,5] depending on the *Spin* and *Spin^c* –structure. In [1], the author defined Seiberg–Witten–like equations on the *Spin* manifold with respect to the generalized self–duality concept and gives them local solutions. Then, in [2] these equations are constructed on the *Spin^c* manifolds and non–trivial local solutions are given to them. Finally, the global solutions of these equations are given on the 8–manifolds endowed with *SU(4)* –structure in [5].

The purpose of this paper is planning to give in two part. One of them is to write down Seiberg–Witten–like equations without using the self–duality concept on the 4 –manifolds and to show similarities with the classical Seiberg–Witten equations. The other one is to define these equations on the 8 –manifolds without using the self–duality concept and to obtain a non–trivial flat solution on the 8 –dimensional Riemannian manifolds. Also, to give them a global solution on the 8 –real–dimensional Kähler manifold for a given negative and constant scalar curvature.

2. Materials and Methods

2.1 *Spin^c* –structure and Dirac operator

Suppose that M is an orientable Riemannian manifold. Hence, there exist an open covering $\{U_\alpha\}_{\alpha \in \mathcal{A}}$ of M with the transitions functions $g_{\alpha\beta}: U_\alpha \cap U_\beta \rightarrow SO(n)$ for TM . If there exists another collection of transition functions $\tilde{g}_{\alpha\beta}: U_\alpha \cap U_\beta \rightarrow Spin^c(n)$ such that the following diagram commutes

$$\begin{array}{ccc}
 & Spin^c(n) & \\
 \tilde{g}_{\alpha\beta} \nearrow & & \downarrow \lambda \\
 U_\alpha \cap U_\beta & \xrightarrow{g_{\alpha\beta}} & SO(n)
 \end{array}$$

That is, $\lambda \circ \tilde{g}_{\alpha\beta} = g_{\alpha\beta}$ and the cocycle condition $\tilde{g}_{\alpha\beta}(x) \circ \tilde{g}_{\beta\gamma}(x) = \tilde{g}_{\alpha\gamma}(x)$ on $U_\alpha \cap U_\beta \cap U_\gamma \neq \emptyset$ is satisfied, then M is called *Spin^c* manifold.

On a *Spin^c* manifold, one can construct $P_{SO(n)}$, $P_{Spin^c(n)}$ and P_{S^1} principal bundles by using principal bundle construction lemma [7]. Also, by using P_{S^1} principal bundle one can construct determinant line bundle

$$\mathcal{L} := P_{Spin^c(n)} \times_l \mathbb{C} = P_{S^1} \times_{U(1)} \mathbb{C}$$

where

$$l_{\alpha\beta} = l: U_\alpha \cap U_\beta \rightarrow Spin^c(n).$$

Moreover, an associated complex vector bundle $\mathbb{S} = P_{Spin^c(n)} \times_{\kappa_n} \Delta_n$ can be constructed by considering spinor representations

$$\kappa_n: Spin^c(n) \rightarrow Aut(\Delta_n)$$

where $\Delta_n = \mathbb{C}^{\frac{n}{2}}$. If the dimension of M is even, then \mathbb{S} spinor bundle splits into two pieces $\mathbb{S} = \mathbb{S}^+ \oplus \mathbb{S}^-$ [4]. The sections of the complex vector bundle are called spinor fields. On the complex vector bundle \mathbb{S} one can define Hermitian inner product as follows:

$$\begin{aligned} \langle, \rangle: \Gamma(\mathbb{S}) \times \Gamma(\mathbb{S}) &\rightarrow \mathbb{C} \\ ([p, \Psi], [p, \Phi]) &\mapsto \langle \Psi, \Phi \rangle = \bar{\Psi} \cdot \Phi. \end{aligned} \quad (2.1)$$

By using Hermitian inner product defined in (2.1) one can associate each spinor Ψ to an endomorphism of \mathbb{S} by the formula

$$\begin{aligned} \Psi\Psi^*: \mathbb{S} &\rightarrow \mathbb{S} \\ \tau &\mapsto \langle \Psi, \tau \rangle \Psi. \end{aligned}$$

Following bundle homomorphisms are useful while studying on spinors. Extended map of κ_n is defined by

$$\kappa: TM \rightarrow \text{End}(\mathbb{S}).$$

Some authors called the map κ a $Spin^c$ -structure on the manifold M [8].

The Clifford multiplication with X is defined

$$X \cdot \Psi := \kappa(X)(\Psi)$$

where $X \in \Gamma(TM)$ and $\Psi \in \Gamma(\mathbb{S})$.

A spinor covariant derivative operator ∇^A is obtained by using an $A: TP_{\mathbb{S}^1} \rightarrow i\mathbb{R}$, $i\mathbb{R}$ -valued 1-form in the principal bundle $P_{\mathbb{S}^1}$ and Levi-Civita connection ∇ on M as follows

$$\nabla^A_X \Psi = d\Psi(X) + \frac{1}{2} \sum_{i < j} \omega_{ij}(X) e_i \cdot e_j(\Psi) + \frac{1}{2} A(X)(\Psi)$$

where $\Psi \in \Gamma(\mathbb{S})$ and $X \in \Gamma(TM)$.

Now we can define the Dirac operator locally as follows.

Definition 1: Let $e = \{e_1, e_2, \dots, e_n\}$ be any local orthonormal frame on $U \subset M$. Then the local expression of the Dirac operator $D_A = \Gamma(\mathbb{S}) \rightarrow \Gamma(\mathbb{S})$ is

$$D_A \Psi = \sum_{i=1}^n e_i \cdot \nabla_{e_i}^A \Psi$$

where $\Psi \in \Gamma(\mathbb{S})$ and $A \in \Omega^1(M, i\mathbb{R})$. Dirac operator decomposes into $D_A = D_A^+ \oplus D_A^-$ in the case of dimension of M is even.

By using κ , another bundle map ρ associated each 2-form to an endomorphism of \mathbb{S} , can be defined on the orthonormal frame $\{e_1, e_2, \dots, e_n\}$ as follows

$$\begin{aligned} \rho: \Lambda^2(T^*M) &\rightarrow \text{End}(\mathbb{S}) \\ \eta = \sum_{i < j} \eta_{ij} e_i \wedge e_j &\mapsto \rho(\eta) = \sum_{i < j} \eta_{ij} \kappa(e_i) \kappa(e_j). \end{aligned}$$

Also ρ can be extended to a complex valued 2-forms [8] such that

$$\rho: \Lambda^2(T^*M) \otimes \mathbb{C} \rightarrow \text{End}(\mathbb{S}).$$

Also ρ can be defined on the half spinor bundles \mathbb{S}^\pm . The half-spinor bundles \mathbb{S}^\pm are invariant under $\rho(\eta)$ for all $\eta \in \Lambda^2(T^*M)$. That is,

$$\begin{aligned} \rho(\eta)(\Psi) &\in \mathbb{S}^+, \forall \Psi \in \mathbb{S}^+ \\ \rho(\eta)(\Psi) &\in \mathbb{S}^-, \forall \Psi \in \mathbb{S}^-. \end{aligned}$$

Then, we obtain the following maps by restriction

$$\begin{aligned} \rho^+(\eta) &= \rho(\eta)|_{\mathbb{S}^+}, \rho^-(\eta) = \rho(\eta)|_{\mathbb{S}^-}. \text{ In this case} \\ \rho^+ &: \Lambda^2(T^*M) \otimes \mathbb{C} \rightarrow \text{End}(\mathbb{S}^+) \end{aligned}$$

is expressed as follows:

$$\rho^+(\eta) = \rho^+(\sum_{i < j} \eta_{ij} e_i \wedge e_j) = \sum_{i < j} \eta_{ij} \kappa(e_i) \kappa(e_j).$$

Note that, the space of $i\mathbb{R}$ -valued 2-forms $\Lambda^2(M, i\mathbb{R})$ is a sub bundle of $\Lambda^2(M, i\mathbb{R}) \times \mathbb{C}$. We consider the sub bundle $W = \rho^+(\Lambda^2(M, i\mathbb{R}))$ of $\text{End}(\mathbb{S})$ to define curvature equation.

In order to be able to give a global solution for the Seiberg-Witten-like equation defined without self-duality, the manifold must be endowed with $SU(4)$ -structure. That guarantees the existence of a Hermitian metric compatible with the complex structure on a Hermitian manifold. On the Hermitian manifold one can construct canonical $Spin^c$ -structure and by using this structure spinorial bundle can be defined with a spinorial connection. Also, Dirac operator is associated with such a connection. As a result Seiberg-Witten-like equation without self-duality can be defined on such manifold and a global solution can be given to it.

In the following, before the global solution is given, a short brief of the Kähler manifolds is given.

2.2 Kähler Manifolds

On the 8-manifolds endowed with $SU(4)$ -structure, there exists an almost complex structure satisfying $J: TM \rightarrow TM, J^2 = -I_d$.

A smooth manifold endowed with an almost complex structure is called an almost complex manifold and denoted by (M, J) .

The almost complex structure J acts on the space of 1-forms as follows:

$$\begin{aligned} J: TM &\rightarrow TM \\ \omega &\mapsto J(\omega)(X) := \omega(JX) \end{aligned}$$

where $\omega \in \Gamma(T^*M)$ and $X \in \Gamma(TM)$. Moreover, J acts on the complexification of the cotangent bundle of M as

$$\begin{aligned} J: T^*M \otimes_{\mathbb{R}} \mathbb{C} &\rightarrow J: T^*M \otimes_{\mathbb{R}} \mathbb{C} \\ \omega \otimes Z &\mapsto J(\omega) \otimes Z. \end{aligned}$$

Since $J^2 = -I_d$, $\pm i$ are eigenvalues of J . Then $T^*M \otimes_{\mathbb{R}} \mathbb{C}$ is the direct sum of

$$T^*M \otimes_{\mathbb{R}} \mathbb{C} = \Lambda^{1,0}(M) \oplus \Lambda^{0,1}(M)$$

where

$$\begin{aligned} \Lambda^{1,0}(M) &= \{Z \in T^*M \otimes_{\mathbb{R}} \mathbb{C} \mid JZ = iZ\} \\ \Lambda^{0,1}(M) &= \{Z \in T^*M \otimes_{\mathbb{R}} \mathbb{C} \mid JZ = -iZ\}. \end{aligned}$$

The space of r -forms is given as:

$$\Lambda^r(M) = \sum_{a+b=r} \Lambda^{a,b}(M)$$

where

$$\Lambda^{a,b}(M) = \text{span}\{x \wedge y \mid x \in \Lambda^a(\Lambda^{1,0}(M)), y \in \Lambda^b(\Lambda^{0,1}(M))\}$$

is the space of (a, b) type complex forms. Finally, Kähler manifold is defined as follows.

Definition 2: Let (M, J) be an almost complex manifold. Then, a Riemannian metric g is called Hermitian metric if it is compatible with the almost complex structure J :

$$g(JX, JY) = g(X, Y)$$

where $X, Y \in \Gamma(TM)$.

The associated smooth 2 –form Φ defined by

$$\Phi(X, Y) = g(X, JY)$$

is called the Kähler 2 –form and satisfies $\Phi(JX, JY) = \Phi(X, Y)$. If Φ is closed then M is called Kähler Manifold and the metric on M is called a Kähler metric.

2.3 Dirac operator on the Kähler Manifolds

In this section, we talk about the canonical $Spin^c$ –structure of a Kähler manifold and its spinor bundle with associated connection. Since the structure group of any Kähler manifold of dimension n is $U(n)$, it admits a canonical $Spin^c$ –structure given by:

$$P_{Spin^c(n)} = P_{U(n)} \times_F Spin^c(2n)$$

where $F: U(n) \rightarrow Spin^c(2n)$ is the lifting map [4]. The associated canonical spinor bundle then has the form:

$$\mathbb{S} \cong \Omega^{(0,*)}(M)$$

where $\Omega^{(0,*)}(M)$ is the direct sum of $\Omega^{(0,1)}(M) \oplus \Omega^{(0,2)}(M) \oplus \dots \oplus \Omega^{(0,i)}(M)$, $i \in \mathbb{N}$. There are two ways to include a spinorial Levi–Civita connection on \mathbb{S} .

The first is obtained by the extension of the connection to forms and the latter is obtained via $Spin^c$ – structure. In this work, we mainly focused on the canonical $Spin^c$ –structure with the following isomorphism:

$$\mathbb{S} \cong \Omega^{(0,*)}(M).$$

On this bundle, we described Dirac operator defined on \mathbb{S} and we give the relation with the Dirac–type operator defined on $\Omega^{(0,*)}(M)$.

In the case of Kähler manifold endowed with a canonical $Spin^c$ –structure, there is a spinorial connection ∇^A on the associated spinor bundle \mathbb{S} induced by an unitary connection 1 –form A on the determinant line bundle \mathcal{L} together with the spinorial Levi–Civita connection ∇ . Also, on the associated spinor bundle one can describe Dirac operator as follows:

Let $\{e_i\}$ $i = 1, \dots, n$ be a local orthonormal frame on M. Then the Dirac operator D^A is given by:

$$D_A \Psi = \sum_{i=1}^n e_i \cdot \nabla_{e_i}^A \Psi.$$

Moreover, by considering Kähler manifolds with $\Omega^{(0,*)}(M)$ associated spinor bundle the Dirac type operator is defined as follows:

Let

$$\bar{\partial}: \Omega^{0,r}(M) \rightarrow \Omega^{0,r+1}(M), \bar{\partial}^*: \Omega^{0,r}(M) \rightarrow \Omega^{0,r-1}(M),$$

given by:

$$\bar{\partial}_0 = \sum_{i=1}^n \bar{Z}_i^* \wedge \nabla_{\bar{Z}_i}, \quad \bar{\partial}_2^* = - \sum_{i=1}^n \iota(\bar{Z}_i)^* \wedge \nabla_{\bar{Z}_i}$$

respectively, where ∇ is the extension of the Levi–Civita connection to $\Omega^{(0,*)}(M)$ and ι is the contraction operator.

Since $\mathbb{S} \cong \Omega^{(0,*)}(M)$, one has

$$D_{A_0} = \sqrt{2} (\bar{\partial}_0 + \bar{\partial}_2^*) \quad (2.3)$$

where A_0 is the Levi–Civita connection of the line bundle $L = \Lambda^2(M)$ of the canonical $Spin^c$ –structure.

2.4 Seiberg–Witten Equations Without Self–Duality on the n – Manifolds

Definition 3: Let (M, g) be a n –dimensional $Spin^c$ manifold. Then Seiberg–Witten Like equations for the pair (A, Ψ) is given by

$$D_A \Psi = 0, \text{ Dirac Equation}$$

$$\rho^+(F_A) = \frac{1}{2} (\Psi \Psi^*)^+, \text{ Curvature Equation} \quad (2.4)$$

where F_A is the curvature of A and $(\Psi \Psi^*)^+$ is the orthogonal projection of $\Psi \Psi^*$ onto $W = \rho^+(\Omega^2(M, i\mathbb{R}))$. In the local orthonormal frame $\{e_1, \dots, e_n\}$,

$$\begin{aligned} (\Psi \Psi^*)^+ &= Proj_W(\Psi \Psi^*) \\ &= \sum_{i < j} \frac{\langle \rho^+(e^i \wedge e^j), \Psi \Psi^* \rangle}{\langle \rho^+(e^i \wedge e^j), \rho^+(e^i \wedge e^j) \rangle} \rho^+(e^i \wedge e^j). \end{aligned}$$

3. Results and Discussion

In this section, we write down the Seiberg–Witten–Like equation on 4 and 8 –dimensional manifolds. Then we compare the solution of these equations with the solution of classical Seiberg–Witten equations on \mathbb{R}^4 [8,9]. Finally, we give a global solution to these equations on 8 –manifolds.

3.1 Seiberg–Witten–like equation on \mathbb{R}^4

In $M = \mathbb{R}^4$ case, the explicit form of the Dirac operator with respect to the $Spin^c(4)$ –structure is given as follows:

$$\begin{aligned} \frac{\partial \psi_1}{\partial x_1} x_4 + A_1 \psi_1 &= i \left(\frac{\partial \psi_1}{\partial x_2} + A_2 \psi_1 \right) + \frac{\partial \psi_2}{\partial x_3} + A_3 \psi_2 \\ &\quad + i \left(\frac{\partial \psi_2}{\partial x_4} + A_4 \psi_2 \right), \\ \frac{\partial \psi_2}{\partial x_1} + A_1 \psi_2 &= -i \left(\frac{\partial \psi_2}{\partial x_2} + A_2 \psi_2 \right) - \frac{\partial \psi_1}{\partial x_3} - A_3 \psi_1 \\ &\quad + i \left(\frac{\partial \psi_1}{\partial x_4} + \frac{1}{2} A_4 \psi_1 \right). \end{aligned}$$

The second equation of the Seiberg–Witten–like equations without self–duality is obtained as follows

$$\begin{aligned} F_{12} + F_{34} &= -\frac{i}{2} (|\psi_1|^2 - |\psi_2|^2), \\ F_{13} - F_{24} &= \frac{1}{2} (\psi_1 \bar{\psi}_2 - \psi_2 \bar{\psi}_1), \\ F_{14} + F_{23} &= -\frac{i}{2} (\psi_1 \bar{\psi}_2 + \psi_2 \bar{\psi}_1). \end{aligned}$$

Notice that, in the case of $M = \mathbb{R}^4$ Seiberg–Witten–like equations without self–duality coincide with the classical Seiberg–Witten equations [8,9].

In the following we define the Seiberg–Witten–like equations on 8 –manifolds.

3.2 Seiberg–Witten– like equation on \mathbb{R}^8

By considering the following $Spin^c$ –structure [5]:

$$\kappa_8: \mathbb{R}^8 \rightarrow \mathbb{C}^{16},$$

$$\kappa_8(e_i) = \begin{bmatrix} 0 & \mu(e_i) \\ -\mu(e_i) & 0 \end{bmatrix} \quad (3.1)$$

where $e_i, i = 1, \dots, 8$ is the standart basis of $\mathbb{R}^8, \mu(e_1) = I_d$ is a 8×8 identity matrix, and for $i = 1, \dots, 8$ explicit form of $\mu(e_i)$ are given by

$$\begin{aligned} \mu(2) &= I_2 \otimes I_2 \otimes m_1, & \mu(3) &= I_2 \otimes I_2 \otimes m_2, \\ \mu(4) &= iI_2 \otimes m_1 \otimes m_1 m_2, & \mu(5) &= iI_2 \otimes m_2 \otimes m_1 m_2, \\ \mu(6) &= -m_1 \otimes m_1 m_2 \otimes m_1 m_2, & \mu(7) &= -m_2 \otimes m_1 m_2 \otimes m_1 m_2, \\ \mu(8) &= -m_1 m_2 \otimes m_1 m_2 \otimes m_1 m_2 \end{aligned}$$

where I_2 is a 2×2 identity matrix and

$$m_1 = \begin{bmatrix} i & 0 \\ 0 & -i \end{bmatrix}, \quad m_2 = \begin{bmatrix} 0 & i \\ i & 0 \end{bmatrix}.$$

By using (3.1) one can obtain explicit form of (2.4) and the solution of these equation can be obtained by the following pair,

$$A = \sum_{i=1}^8 -2ix_i dx^i$$

and

$$\Psi = \left(0, 0, 0, e^{\sum_{j=1}^8 -\frac{1}{2}x_j^2}, 0, 0, e^{\sum_{j=1}^8 -\frac{1}{2}x_j^2}, 0\right).$$

Here (A, Ψ) is the local, non-trivial but flat (*ie.* $F_A = 0$) solution of the Seiberg–Witten–like equation without self-duality with respect to $M = \mathbb{R}^8$.

In the next subsection, a global solution to the Seiberg–Witten–like equations without self-duality is given on 8–real–dimensional Kähler Manifolds.

3.3 Seiberg–Witten–Like Equations on the 8–Real–Dimensional Kähler Manifold

Let (M, g, J) be a 8–real–dimensional Kähler manifold endowed with a canonical $Spin^c$ –structure and $e_1, e_2 = J(e_1), e_3, e_4 = J(e_3), e_5, e_6 = J(e_5), e_7, e_8 = J(e_7)$, be a local orthonormal frame with the dual basis $\{e_1, e_2, e_3, e_4, e_5, e_6, e_7, e_8\}$. Then the Kähler 2–form has the form

$$\Phi = e_1 \wedge e_2 + e_3 \wedge e_4 + e_5 \wedge e_6 + e_7 \wedge e_8.$$

Under the action Φ , one gets the following decomposition

$$\mathbb{S} = \mathbb{S}_0 \oplus \mathbb{S}_1 \oplus \mathbb{S}_2 \oplus \mathbb{S}_3 \oplus \mathbb{S}_4,$$

where

$$\begin{aligned} \mathbb{S}_0 &= \{\Psi \in \mathbb{S} \mid \Phi \Psi = 4i\Psi\}, \\ \mathbb{S}_1 &= \{\Psi \in \mathbb{S} \mid \Phi \Psi = 2i\Psi\}, \\ \mathbb{S}_2 &= \{\Psi \in \mathbb{S} \mid \Phi \Psi = 0\}, \\ \mathbb{S}_3 &= \{\Psi \in \mathbb{S} \mid \Phi \Psi = -2i\Psi\}, \\ \mathbb{S}_4 &= \{\Psi \in \mathbb{S} \mid \Phi \Psi = -4i\Psi\}. \end{aligned}$$

Accordingly, $f: i^4 e_1 \cdot e_2 \cdot e_3 \cdot e_4 \cdot e_5 \cdot e_6 \cdot e_7 \cdot e_8: \mathbb{S} \rightarrow \mathbb{S}$ endomorphism, the complex spinor bundle \mathbb{S} splits into $\mathbb{S} = \mathbb{S}^+ \oplus \mathbb{S}^-$

where

$$\begin{aligned} \mathbb{S}^+ &= \mathbb{S}_0 \oplus \mathbb{S}_2 \oplus \mathbb{S}_4 \cong \Lambda^{0,4}(M) \oplus \Lambda^{0,2}(M) \oplus \Lambda^{0,0}(M), \\ \mathbb{S}^- &= \mathbb{S}_1 \oplus \mathbb{S}_3 \cong \Lambda^{0,3}(M) \oplus \Lambda^{0,1}(M). \end{aligned}$$

Let Ψ_0 be a spinor in $\mathbb{S}_4 \cong \Omega^{0,0}(M)$ corresponding to constant function 1 in the chosen coordinates

$$\Psi_0 = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 1 \end{bmatrix}. \text{ By using } \Psi_0, \text{ one has}$$

$$\frac{(\Psi_0 \Psi_0^*)^+}{2} = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \frac{-1}{4} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \frac{1}{4} \end{bmatrix}.$$

Let A_0 be the connection on the S^1 –principal bundle P_{S^1} induced by means of the Levi–Civita connection ∇ in the line bundle $L = \Omega^{0,2}(M)$ of the canonical $Spin^c$ –structure [4]. Accordingly, the corresponding Dirac operator $D_{A_0}: \Gamma(\mathbb{S}^+) \rightarrow \Gamma(\mathbb{S}^-)$ coincides with $\sqrt{2}(\partial_0 \oplus \bar{\partial}_2^*)$. Also, the curvature of the connection 1–form A_0 is given by

$$F_{A_0} = i\rho_{ric} \quad (3.2)$$

where $\rho(X, Y) = (X, Y) = g(X, J \circ Ric(Y))$ and $Ric: TM \rightarrow TM$ denotes the Ricci tensor. Since the almost complex structure J and the Ricci tensor Ric commute, one has

$$\begin{aligned} \rho_{ric} &= -R_{11}e_1 \wedge e_2 - R_{33}e_3 \wedge e_4 - R_{13}(e_1 \wedge e_4 - e_2 \wedge e_3) \\ &+ R_{14}(e_1 \wedge e_3 - e_2 \wedge e_4) - R_{15}(e_1 \wedge e_6 - e_2 \wedge e_5) \\ &+ R_{16}(e_1 \wedge e_5 + e_2 \wedge e_6) - R_{17}(e_1 \wedge e_8 - e_2 \wedge e_7) \\ &+ R_{18}(e_1 \wedge e_7 + e_2 \wedge e_8) - R_{35}(e_3 \wedge e_6 - e_4 \wedge e_5) \\ &+ R_{36}(e_3 \wedge e_5 + e_4 \wedge e_6) - R_{37}(e_3 \wedge e_8 - e_4 \wedge e_7) \\ &+ R_{38}(e_3 \wedge e_7 + e_4 \wedge e_8) - R_{55}e_5 \wedge e_6 - R_{77}e_7 \wedge e_8 \\ &- R_{57}(e_5 \wedge e_8 - e_6 \wedge e_7) + R_{58}(e_5 \wedge e_7 - e_6 \wedge e_8). \end{aligned}$$

In the following a global solution is given for the appropriate Ricci tensor.

Theorem 1.

Let (M, g, J) be an 8–real–dimensional Kähler manifold. Then for a given negative and constant scalar curvature s ($A_0, \Psi = \sqrt{-2s}\Psi_0$) is the solution of the Seiberg–Witten–like equations without self-duality.

Proof. Since $\Psi = \sqrt{-2s}\Psi_0 \in \Omega^{0,0}(M)$ and Ψ is the spinor field corresponding to the constant function 1, by using (2.3), one gets $D_{A_0} \Psi \equiv 0$. Satisfying the curvature equation remains. To achieve this, Ric must be taken as follows:

$$Ric = [R_{ij}]_{8 \times 8} = \begin{cases} \frac{s}{8} & i = j \\ 0 & i \neq j, \end{cases}$$

where s is the negative and constant. By using Ric in (3.2), one gets $\rho^+(F_{A_0}) = i\rho^+(\rho_{Ric})$ which means $\rho^+(F_{A_0}) = \frac{(\Psi\Psi^*)^+}{2}$.

4. Conclusion

We give a global solution to the Seiberg–Witten–like equations on 8–real–dimensional Kähler manifolds for a given negative and constant scalar curvature.

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